

1 Satellite-based Analysis of Ocean-Surface Stress across the Ice-free 2 and Ice-covered Polar Oceans

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6 **Abstract.** Ocean-surface stress is a critical driver of polar sea ice dynamics, air-sea interactions, and ocean circulation. This
7 work provides a daily analysis of ocean-surface stress on 25-km Equal-Area Scalable Earth (EASE) Grids across the ice-free
8 and ice-covered regions of the polar oceans (2011-2021 for Arctic, 2013-2021 for Antarctic), covering latitudes north of 60°N
9 in the Arctic and south of 50°S in the Antarctic and Southern Ocean. Ocean-surface stress is calculated using a bulk
10 parameterization approach that combines ocean-surface winds, ice motion vectors, and sea surface height (SSH) data from
11 multiple satellite platforms. The analysis captures significant spatial and temporal variability in ocean-surface wind stress and
12 the resultant wind-driven Ekman transport, while providing enhanced spatiotemporal resolution. Two sensitivity analyses are
13 conducted to address key sources of uncertainty. The first addresses the fine-scale variability in SSH fields, which was
14 mitigated using a 150-km Gaussian filter to smooth three-day SSH datasets and enhance compatibility with the other monthly
15 product, followed by linear interpolation to achieve daily resolution. The second investigates uncertainty in the ice-water drag
16 coefficient, which revealed that variations in the coefficient have a proportional influence on the computed ocean-surface stress
17 under the tested conditions. These uncertainties are most pronounced during winter, with median values reaching 20% in the
18 Arctic and 40% in the Southern Ocean. Validation efforts using Ice-Tethered Profiler velocity records revealed weak to
19 moderate correlations with satellite-derived stress ($r = 0.4-0.8$) between observed surface velocities and satellite-derived
20 estimates (Ekman + geostrophic) at daily resolution, with significantly improved agreement when averaged to weekly means.
21 This dataset is publicly available at <https://doi.org/10.5281/zenodo.15534576> (Liu & Yu, 2024).

22

23 1 Introduction

24 Earth's polar regions have undergone profound changes over the past decades, with sea ice playing a central role in the polar
25 climate system. By modulating heat, momentum, and freshwater exchanges at the atmosphere-ice-ocean boundary, sea ice
26 directly influences global climate dynamics (Meehl, 1984; Stammerjohn et al., 2012). In the Arctic, rapid sea ice decline has
27 transitioned the region from predominantly thick, multiyear ice to thinner, more dynamic ice, with increased interannual
28 variability (Comiso et al., 2008; Stroeve and Notz, 2018; Moore et al., 2022; Babb et al., 2022). Meanwhile, Antarctic sea ice
29 trends have shown greater complexity, with a modest long-term increase observed until the mid-2010s, followed by a record
30 loss in 2017 and a subsequent continued decline (Liu et al., 2004; Parkinson, 2019; Turner et al., 2022; Purich & Doddridge,
31 2023). These changes in sea ice extent and thickness have significant implications for polar systems and global climate
32 feedbacks, influencing the Arctic's ability to regulate planetary heat, as well as impacting marine ecosystem, carbon cycling,
33 nutrient distribution, and thermohaline circulation (Talley, 2013; Campbell et al., 2019).

34 Atmospheric circulation is a primary driver of sea ice dynamics and variability. Geostrophic winds, for instance, account for
35 over 70% of sea ice velocity variability (Thorndike and Colony, 1982; Maeda et al., 2020), while broader climate modes,
36 including the Arctic Oscillation, Pacific Decadal Oscillation, and Southern Annular Mode, influence ice extent and distribution
37 (Rigor et al., 2002; Park et al., 2018; Lefebvre et al., 2004). These wind-driven processes interact with sea ice to modify ocean-
38 surface stress, impacting Ekman dynamics and the transport of heat, salt, and nutrients (Yang, 2006, 2009; Meneghello et al.,
39 2018). This feedback mechanism, often described as the "ice-ocean governor" (Meneghello et al., 2017), plays an important
40 role in regulating polar freshwater storage and circulation (Marshall and Speer, 2012; Abernathey et al., 2016; Ma et al., 2017).
41 Surface stress plays a pivotal role in driving Arctic Ocean circulation by mediating the transfer of momentum from the
42 atmosphere to the ocean. In the Arctic, sea ice acts as a modulator of this momentum exchange, either dampening or amplifying
43 the transfer depending on its concentration and mechanical properties. Recent projections indicate that as the Arctic climate
44 warms, sea ice will become thinner and less extensive, leading to a more efficient transfer of wind energy to the ocean surface
45 (Mulwijk et al., 2024). This enhanced momentum transfer is expected to accelerate surface currents, increase ocean kinetic
46 energy, and intensify vertical mixing processes (Martin et al. 2014; Martin et al. 2016). However, current climate models
47 exhibit considerable uncertainty in simulating these processes due to simplified representations of atmosphere-ice-ocean
48 interactions. Therefore, developing observationally based surface stress products is essential for validating and improving
49 model simulations, leading to more accurate predictions of future Arctic Ocean dynamics and their global implications.

50 To address the complexity of ice-ocean interactions, recent modeling advances have highlighted the pivotal role of sea ice
51 form drag in governing momentum exchange at the ocean-ice-atmosphere interface. Tsamados et al. (2014) introduced a
52 physically grounded parameterization of ice form drag that accounts for ice morphological features—such as ridges, floe edges,
53 and melt pond geometry—and demonstrated that spatial and temporal variability in drag coefficients can substantially
54 influence sea ice dynamics and the spin-up of the Arctic Ocean. Extending this approach, Sterlin et al. (2023) implemented a
55 variable ice form drag scheme in the NEMO-LIM3 ocean-sea ice model and found that it exerts a pronounced control over

56 ocean surface stress patterns, mixed layer depth, sea surface salinity, and upper ocean temperature across both polar regions.
57 These modeling efforts reveal that ice form drag is not merely a secondary detail, but a first-order process in polar ocean
58 circulation and surface forcing. However, the representation of ice–ocean drag—often quantified through the coefficient—
59 remains highly uncertain, as it can vary markedly with environmental conditions including ice concentration, surface roughness,
60 and the presence of waves (Lüpkes & Gryanik, 2015; Brenner et al., 2021). This highlights the growing need for
61 observationally based estimates of ocean surface stress that can support parameterization efforts, constrain model behavior,
62 and improve the physical realism of coupled ocean–ice simulations.

63 Despite significant advancements in understanding these processes, direct measurements at the ice-ocean interface remain
64 limited, with most data concentrated in the Arctic’s Canada Basin (Smith et al., 2019; Regan et al., 2019). Satellite remote
65 sensing has been instrumental in addressing these gaps, providing open ocean-surface wind retrievals available since 1988 (Yu
66 & Jin, 2014a) and tracking sea ice motions since 1978 (Cavalieri et al., 1996). Recent advances in satellite altimetry further
67 enable high-resolution monitoring of sea surface height (SSH) changes, offering new insights into mesoscale ocean dynamics
68 (Armitage et al., 2016, 2017; Prandi et al., 2021).

69 Building upon the concepts developed in previous studies (Yang, 2006, 2009; Meneghello et al., 2018), this analysis utilizes
70 recent satellite-based datasets on wind, ice motion, and SSH to analyze ocean-surface stress across both ice-free and ice-
71 covered polar seas. Specifically, we present a daily analysis of ocean-surface stress at 25-km resolution using Equal-Area
72 Scalable Earth (EASE, see glossary in Table A1 for more details) Grids from 2011 to 2021 for Arctic and 2013-2021 for
73 Antarctic, covering latitudes north of 60°N in the Arctic and south of 50°S in the Antarctic and Southern Ocean (Figure 1).
74 Section 2 provides a description of the satellite datasets used and processing steps, along with the methods for calculating
75 ocean-surface stress and Ekman circulation. Section 3 presents the time-mean patterns and variability of the derived surface
76 stress and Ekman pumping fields. Section 4 addresses quantification of uncertainties in the analysis, including sensitivity to
77 the ice-water drag coefficient and comparisons of with in-situ data.

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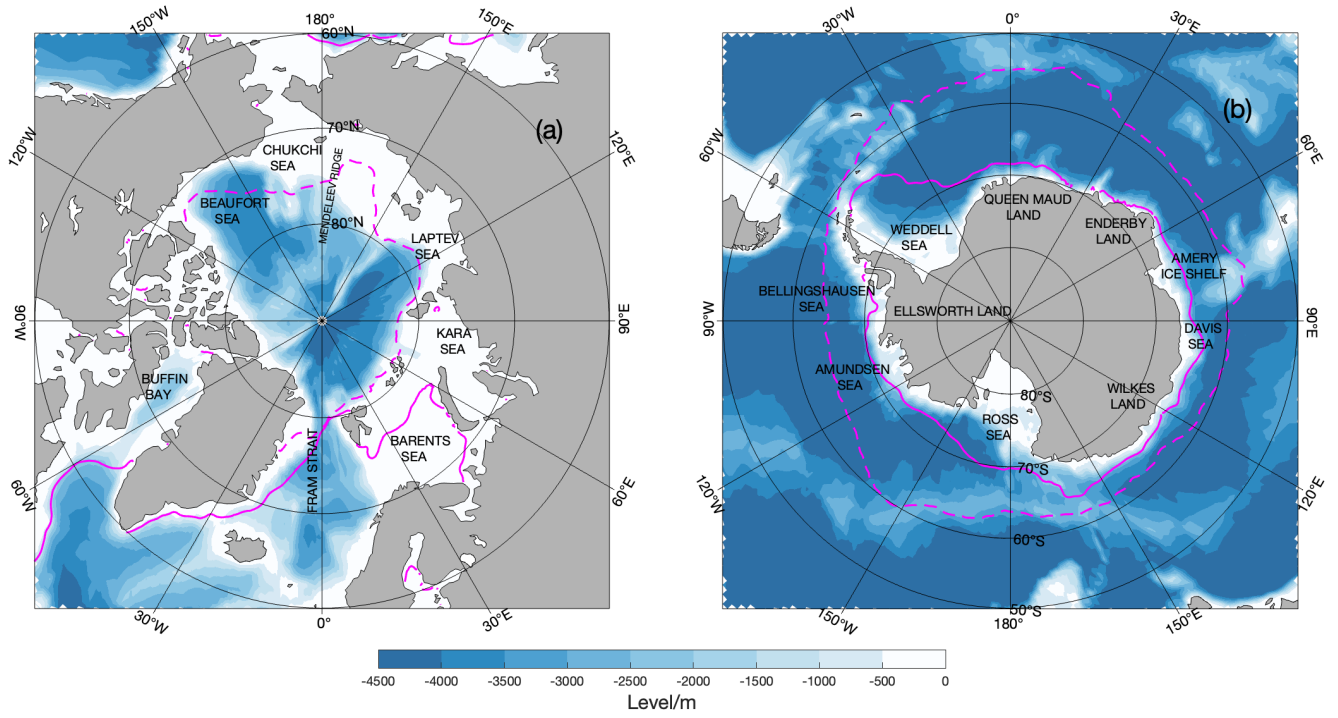


Figure 1: Study region in (a) Arctic and (b) Southern Oceans. Blue shading represents the bathymetry in meter. Solid and dashed magenta lines indicate the median sea ice extent boundaries for March and September, respectively, defined by areas with sea ice concentration.

2 Data, Method and Processing of the Analysis

2.1 Calculation of Ocean-Surface Stress and the Ekman Transport

The ocean-surface stress is estimated using the methodology proposed by Yang (2006, 2009), with modifications by Meneghello et al. (2018). The total ocean-surface stress (τ_o) is calculated as a weighted linear combination of ice–water stress (τ_{iw}) and air–water stress (τ_{aw}), based on the fractional sea ice concentration:

$$\tau_o = \alpha \tau_{iw} + (1 - \alpha) \tau_{aw} \quad (1)$$

where α is set to 0 for the ice-free surfaces (defined as sea ice concentration less than 15%) and 1 for ice-covered surfaces (defined as sea ice concentration exceeding 15%). The stresses τ_{iw} and τ_{aw} are parameterized using quadratic drag laws:

$$\tau_{iw} = \rho_w C_{D,iw} |U_{ice} - U_e - U_g| (U_{ice} - U_e - U_g) \quad (2)$$

and

$$\tau_{aw} = \rho_a C_{D,aw} |U_{10}| U_{10} \quad (3)$$

where U_{ice} , U_e , U_g , and U_{10} are the local ice motion, Ekman velocity, geostrophic velocity, and equivalent neutral wind at 10-m height, respectively. $\rho_w = 1027.5 \text{ kg m}^{-3}$ and ρ_a represent the densities of water and air. In this product, τ_{aw} is taken directly from existing satellite wind products (Yu and Jin 2014a, b).

98 $C_{D,iw}$ is the ice-water drag coefficient and $C_{D,iw} = 5.5 \times 10^{-3}$ is adopted in this product as it is a commonly recognized value.
99 It is worth noting that, due to the limited availability of direct observations, $C_{D,iw}$ is identified as a key source of uncertainty.
100 A sensitivity analysis is therefore provided in the following section to evaluate its potential impact.
101 In Equation (2), surface ocean velocity expressed as the sum of U_g and U_e . The representation of ocean surface stress is known
102 to be highly sensitive to the assumed surface velocity used in the drag formulation. A range of approaches has been employed
103 in past studies—incorporating U_e , U_g , or even assuming zero ocean motion—each with markedly different implications. For
104 instance, Zhong et al. (2018) showed that mean Ekman pumping in the Beaufort Sea can differ by over 50% depending on the
105 inclusion of geostrophic flow. Wu et al. (2021) reported similar sensitivities in the Nordic Seas, while earlier works by Zhong
106 et al. (2015) and Ma et al. (2017) further detailed the variability across Arctic regimes. As a result, stress-based diagnostics
107 remain sensitive to parameterization choices, and conclusions should be interpreted with that uncertainty in mind.
108 The geostrophic velocity U_g can be calculated from dynamic ocean topography datasets (McPhee 2013; Armitage et al. 2016,
109 2017). The Ekman velocity U_e , which moves at an angle of 45° to the right of the ocean-surface stress in the Northern
110 Hemisphere, is calculated as:

$$111 \quad U_e = \frac{\sqrt{2}e^{-i(\pi/4)}}{f\rho_w D_e} \tau_o \quad (4)$$

112 where f is the Coriolis parameter, and D_e is the Ekman layer depth (20 m, Meneghello et al., 2018). Since U_e and τ_o are
113 interdependent in Eqs. (1) and (4), a modified Richardson iteration method is applied to solve them iteratively until converge
114 is achieved, starting with $U_e = 0$ in the first iteration (Yang 2006).

115 Subsequently, the vertical Ekman velocity w_e can be calculated as follows:

$$116 \quad w_e = \frac{1}{f\rho_w} \nabla \times \tau_o \quad (5)$$

117 A positive w_e indicates upwelling, while a negative w_e corresponds to downwelling.

118

119 2.2 Data Description

120 The calculation of total ocean-surface stress (Eqs. (1)–(4)) requires the following input datasets: ocean-surface wind stress
121 (τ_{aw}), sea ice concentration (α), sea ice motion (U_{ice}), and dynamic topography for geostrophic velocity (U_g). A brief
122 description of each satellite-based dataset is given in Table 1.

123

124 **Table 1: Gridded satellite datasets used in the work.**

Variable	Source	Resolution	Period	Reference
Surface	OAFlux2	Daily, 0.25°	1988-present	Yu & Jin, 2014a, 2014b
Wind Stress τ_{aw}				
Ice Motion U_{ice}	Polar Pathfinder v4	Daily, 25 km	1978-2023	Tschudi et al., 2019

Geostrophic U_g	multi-altimeter dataset	3-Day, 25 km	2011-2021 (Arctic) 2013-2021 (Antarctic)	Prandi et al., 2021
Sea Ice Concentration α	NSIDC0051, v2	Daily, 25 km	1988-present	DiGirolamo et al., 2022

In this product, the air-water wind stress is taken from OAFlux2 (Yu & Jin, 2014a, 2014b), a satellite-derived 0.25-degree gridded air-sea flux daily analysis (1988 to present) developed under the auspices of NASA's Making Earth System Data Records for Use (MEaSUREs) program (Yu, 2019). OAFlux2 winds are synthesized from 19 active and passive satellite wind sensors and wind stress are calculated from the Coupled Ocean-Atmosphere Response Experiment (COARE) bulk algorithm version 3.6 (Fairall et al., 2003).

Daily sea ice motion vectors for the Arctic and Antarctic regions are obtained from the National Snow and Ice Data Center's (NSIDC) Polar Pathfinder Daily 25 km EASE-Grid Sea Ice Motion Vectors, Version 4 (Tschudi et al., 2019, 2020), covering the period from 1978 through 2023. The ice motion fields are derived from multiple sources, including passive microwave radiometers (e.g., SSM/I, AMSR-E), visible and infrared sensors (e.g., AVHRR, MODIS), scatterometers (e.g., QuikSCAT), drifting buoys (e.g., IABP), and atmospheric reanalysis winds. Feature-tracking algorithms are applied to sequential satellite images to identify ice displacement, while optimal interpolation techniques combine the various data sources to produce daily sea ice motion estimates. The resulting vectors represent sea ice displacement over a 24-hour period and are gridded onto a 25 km EASE grid 2.0 (EASE2).

Geostrophic velocity in the Arctic and Antarctic are obtained from the CLS/PML multi-altimeter combined Arctic/Antarctic Ocean sea level dataset (Prandi et al., 2021). This dataset spans latitudes north of 50°N on a 25 km EASE2, with a temporal resolution of one grid point every 3 days. Covering the Arctic from 2011 to 2021 and the Antarctic from 2013 to 2021, the CLS dataset mitigates the spurious meridional signals often introduced by the longer sampling intervals of CryoSat-2 observations (Auger et al., 2022).

The Sea Ice Concentrations from Nimbus-7 SMMR and DMSP SSM/I–SSMIS Passive Microwave Data, Version 2 (NSIDC-0051, Cavalieri et al., 1996; DiGirolamo et al., 2022) is used to define the daily ice boundary based on the 15% ice concentration threshold. NSIDC-0051 provides a reliable, long-term record of sea ice concentration, making it valuable for studying sea ice conditions and large-scale climate variability (Parkinson, 2019). Widely recognized for its accuracy, the dataset is frequently used to validate and improve climate model simulations. The daily dataset is available from 1987 to the present and provide a coverage on a 25 km resolution polar stereographic grid for the both polar regions.

The analysis period ends in 2021 to maintain consistency with the most reliable iteration of the ongoing refinement of the associated satellite products as listed in Table 1. While this choice limits the temporal extent, the framework itself remains

flexible and can be readily extended as newer, better-resolved datasets become available. With all considered, the study period for Antarctica is constrained to six years (2013–2021), while an eight-year period (2011–2021) is maintained for the Arctic.

2.3 Data processing procedure

Using the methodology described in Eqs. (1)–(5) and the input data listed in Table1, the workflow for processing and analysing data to calculate ocean-surface stress and derive vertical Ekman velocity is shown in Figure 2.

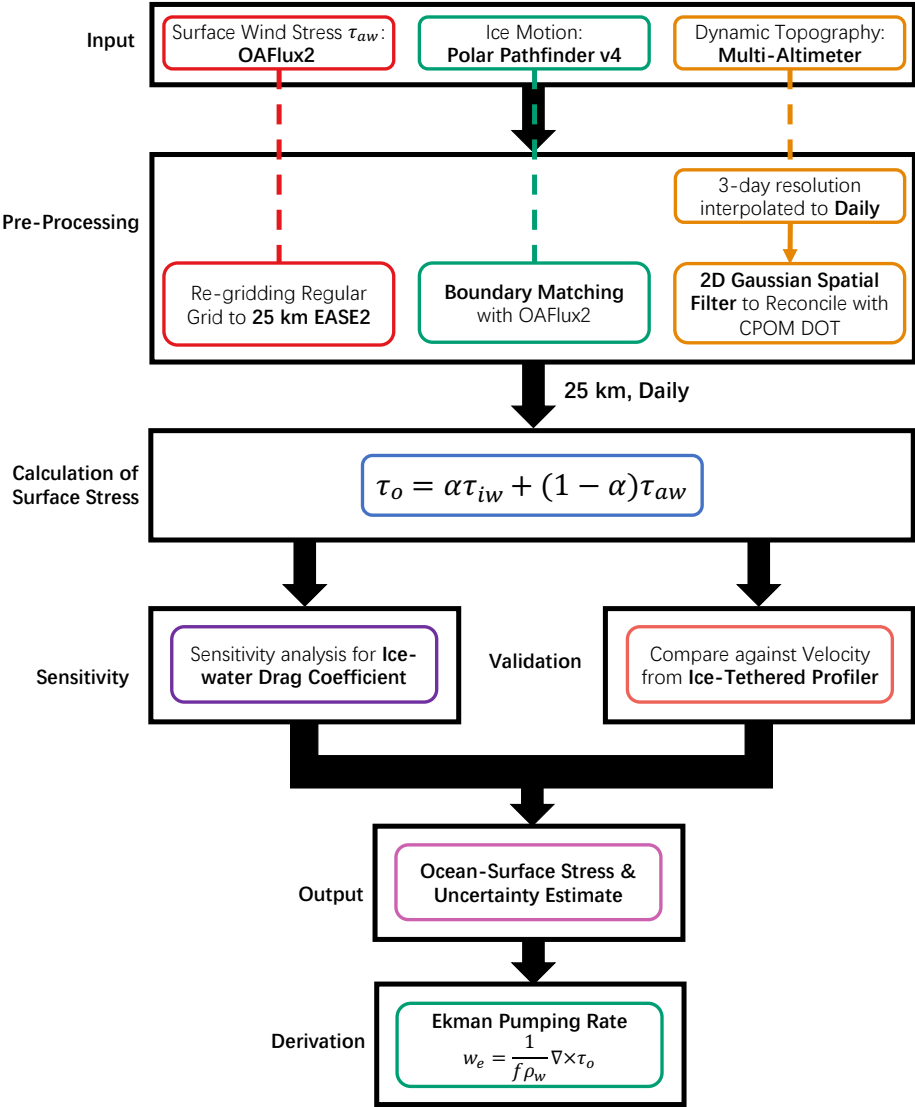


Figure 2: Workflow for data processing and analysis to calculate ocean-surface stress and derive vertical Ekman velocity.

159 All datasets are interpolated onto a common 25 km EASE grid format, providing uniform spatial resolution and facilitating
 160 consistent analysis across the Arctic and Antarctic regions. Noting that the 25 km resolution may introduce uncertainties near
 161 the 15% sea ice concentration boundary, as such coarse resolution can obscure sharp gradients in the marginal ice zone and
 162 misclassify mixed ice–water grid cells (e.g., Meier, 2005; Ivanova et al., 2015).
 163 Despite efforts to merge ice-edge boundaries across multiple satellite products, we note that data gaps and inconsistencies
 164 become more pronounced after 2019. This is primarily due to increasing divergence between input datasets used to define sea
 165 ice concentration and motion. As a result, the quality of the derived surface stress fields may be reduced, particularly in the
 166 marginal ice zone (MIZ), where small changes in ice coverage can significantly affect stress partitioning. Users should exercise
 167 caution when applying this dataset to study MIZ dynamics after 2019, and we recommend validating results against
 168 independent sources where possible.
 169 Temporal sampling frequency plays a critical role in determining the accuracy and interpretability of ocean surface stress
 170 estimates. Daily and sub-daily sampling is often needed to capture short-term variability in wind and sea ice motion, which
 171 directly affects transient stress fluctuations and high-frequency Ekman responses (Meneghello et al., 2018; Regan et al., 2020).
 172 Conversely, monthly-averaged fields may overly smooth dynamic features and obscure important stress events, particularly in
 173 regions with strong synoptic variability. In this context, the 3-day sampling of the CLS/PML altimetry product offers a useful
 174 compromise: it resolves large-scale mesoscale dynamics more consistently than monthly data while maintaining better signal-
 175 to-noise properties than noisy daily sea surface height reconstructions in ice-covered regions (Prandi et al., 2021). Given the
 176 limitations of satellite altimetry in the polar oceans, the CLS dataset provides a crucial improvement by reducing spurious
 177 meridional errors and enabling more consistent estimation of geostrophic velocities and their role in modulating surface stress
 178 (Auger et al., 2022). We revisit the implications of time-averaging choices for surface stress/derived velocity fields in Section
 179 4.2.
 180 While the CLS/PML product offers improved temporal resolution and geophysical realism at finer spatial scales, its use of
 181 multiple altimeters and interpolation techniques can introduce high-frequency structures that remain difficult to validate given
 182 the sparse in situ coverage at high latitudes (Prandi et al., 2021; Auger et al., 2022). In contrast, the CPOM Dynamic Ocean
 183 Topography (DOT) dataset (2003-2021) from the Centre for Polar Observation and Modelling (CPOM; Armitage et al., 2016,
 184 2017), despite its coarser resolution and monthly cadence, has seen wider adoption and validation across climate-scale Arctic
 185 studies (e.g., Meneghello et al., 2018; Zhong et al., 2018; Lin et al., 2023), making it a valuable benchmark for cross-
 186 comparison. To reconcile the strengths of both datasets, we apply a two-dimensional Gaussian spatial filter to smooth CLS
 187 fields, aligning their effective resolution with CPOM and improving interpretability of large-scale patterns. This hybrid
 188 approach leverages the temporal detail of CLS while benefiting from the broader-scale reliability of CPOM, offering a more
 189 balanced foundation for stress estimation and error characterization in polar oceanographic applications.
 190 We employed a 2D Gaussian filter with a standard deviation of 75 km to improve consistency and interpretability between
 191 CLS and CPOM DOT datasets, which have different resolutions and small-scale characteristics. A sensitivity test is conducted
 192 to determine the optimal filter radius, ranging from 50 km to 250 km. Smaller filters (e.g., <50 km) preserve small-scale

193 variability but may complicate the interpretation of large-scale features, while larger filters (e.g., >250 km) can excessively
 194 smooth mesoscale processes, such as boundary currents, reducing the dataset's ability to capture key processes of polar
 195 dynamics.

196 To find the optimal filter size, a series of tests were conducted for 2011. The effectiveness of each filter setting is evaluated
 197 using the Root Mean Squared Deviation (RMSD):

$$198 \quad RMSD = \sqrt{\frac{1}{N} \sum_{i,j} (w_{e\ i,j} - w_{eref,(i,j)})^2} \quad (6)$$

199 where w_e represents the local vertical Ekman velocity w_e derived from the CLS dataset, filtered with a specific Gaussian filter
 200 size (e.g., 100 km, 150 km, etc.), and w_{eref} is the reference vertical Ekman velocity calculated using the CPOM dataset. N is
 201 the total number of the grid points with sea ice coverage.

202 The unfiltered CLS dataset exhibits clear seasonal variations in RMSD, with values peaking at 25 cm/day during winter and
 203 decreasing in summer (Figure 3a). Applying a Gaussian filter significantly enhances agreement with the CPOM dataset,
 204 reducing RMSD by 10–15 cm/day for most of the year. However, in late summer the reduction is only 2–5 cm/day.

205 Increasing the filter size further enhances spatial agreement (Figure 3b). From no-filter to a 100 km filter, the annual mean
 206 RMSD is reduced to 17 cm/day, and increasing the filter size to 150 km further lowers the RMSD to 15.5 cm/day. The standard
 207 deviation of daily RMSD is also reduced by half with 150 km filter compared to the unfiltered results. However, larger filter
 208 sizes (e.g. 200 km and 250 km) yield only marginal additional improvements. Therefore, the 150 km Gaussian filter is selected
 209 as a practical and effective balance between preserving spatial features and minimizing small-scale variability for this work.

210 Figure 3c–h demonstrates the impact of varying filter sizes on the spatial structures of τ_o and w_e on March 15, 2011. Without
 211 filtering, the CLS dataset exhibits residual meridional striping due to satellite sampling artifacts (Auger et al., 2022). This
 212 pattern is significantly suppressed with a 150 km Gaussian filter. Between the filtered CLS-derived w_e (150 km) and CPOM-
 213 derived w_e , the correlation coefficients improving markedly from 0.77 (no filter) to over 0.95 ($p < 0.05$).

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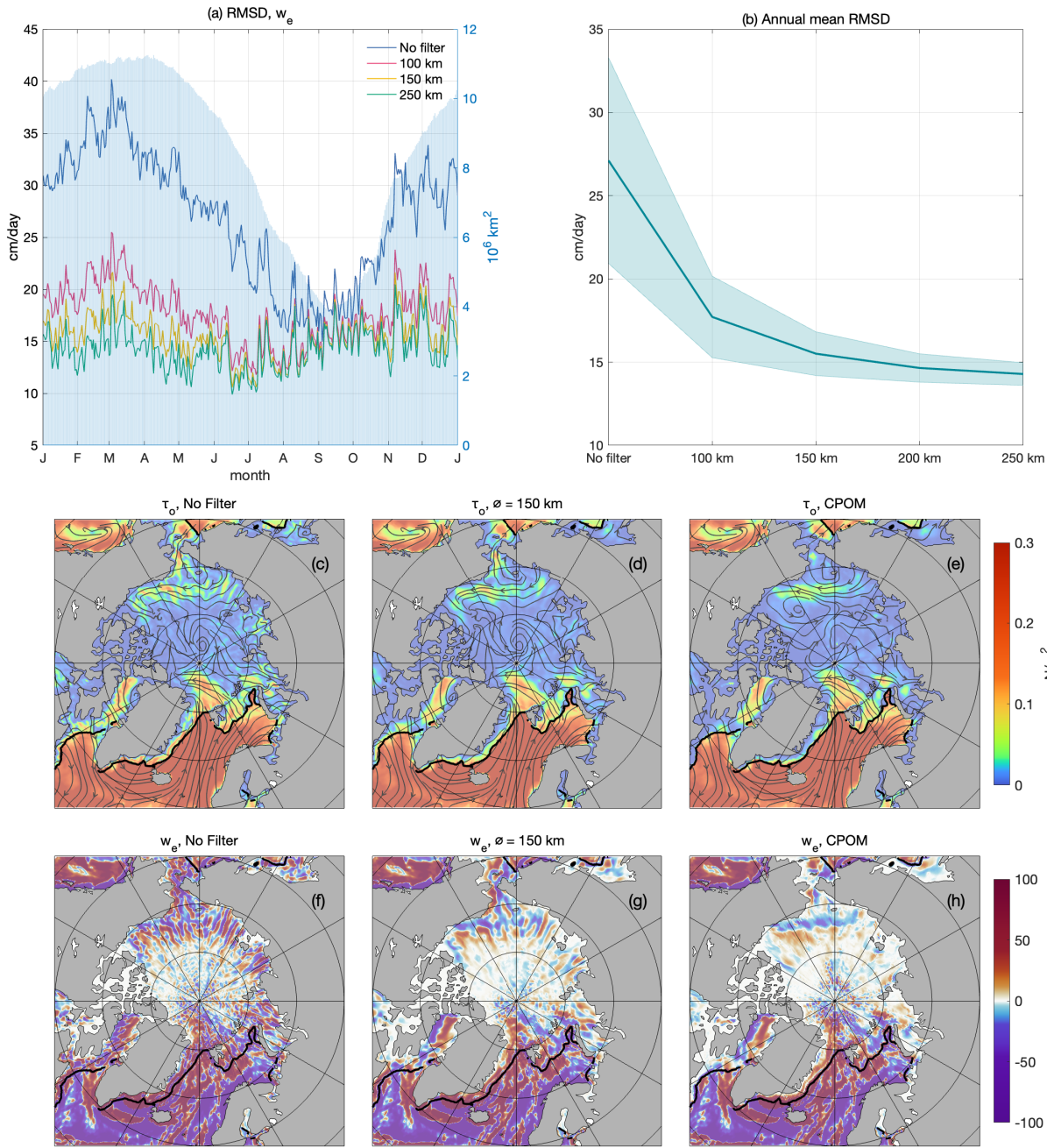


Figure 3: Area-averaged ocean-surface stress τ_o and vertical Ekman velocity w_e regarding Gaussian filter setting. (a) Annual cycle of root mean squared deviation (RMSD) of w_e over 2011. Blue shades show total ice-cover areas (right axis). (b) Annual mean RMSD of w_e with shading indicating one standard deviation over a year. (c) Snapshot of τ_o with unfiltered CLS (3/15/2011). (d) Same as c but with 150 km Gaussian filter. (e) Same as c but with CPOM. (f-h) Same as (c-e) but for w_e . Streamlines in (c-e) show the direction of τ_o . Black contours in (c-h) mark the 15% ice concentration on 3/15/2011.

222 3 Results and Regional Statistics

223 3.1 Arctic Ocean

224 In this section, we provide a concise overview of the surface stress and the corresponding Ekman velocity fields. Figure 4a
225 shows the time-averaged ocean-surface stress (τ_o) field across the Arctic for 2011–2021. The highest τ_o appears in the ice-
226 free Nordic Seas, where strong wind-ocean interactions drive surface stress exceeding 0.3 N/m². In contrast, sea ice reduces
227 momentum transfer and lower the τ_o in ice-covered regions. In the Seasonal Ice Zone (SIZ), marked by the March and
228 September sea ice boundaries, τ_o typically remains below 0.05 N/m². Within the Perennial Ice Zone (PIZ), bounded by the
229 September sea ice boundary, it drops further to below 0.02 N/m².

230 The seasonal cycle of τ_o is the dominant temporal variability across the Arctic (Stroeve and Notz, 2018). The standard
231 deviation (STD) shows a spatial distribution similar to the time-averaged τ_o (Figure 4b), with high variability (>0.1 N/m²) in
232 ice-free regions like the Nordic Seas. Variability is significantly suppressed in the SIZ and PIZ, with values below 0.02 N/m²
233 and 0.01 N/m², respectively. The coefficient of determination (R^2 , here is calculated as the proportion of variance explained
234 by the seasonal cycle, i.e., $R^2 = 1 - \frac{\sum_i (y_i - y_{seasonal})^2}{\sum_i (y_i - \bar{y})^2}$) shows that in open-ocean regions, 40–60% of variance is explained by
235 seasonal variability (Figure 4c). In ice-covered areas, this ratio drops to less than 30%.

236 The time-mean Ekman pumping rate (w_e), alongside its STD and R^2 patterns are given in Figure 4d-f. Strong upwelling (>50
237 cm/day) is observed in the Nordic Seas, while strong downwelling (<-10 cm/day) occurs in the Beaufort and Chukchi Seas.
238 The spatial pattern of STD w_e is similar to that of τ_o (Figure 4e). Seasonal variability ranges from 10–20 cm/day in ice-free
239 regions, 4–6 cm/day in SIZ, and falls below 4 cm/day in the PIZ. Seasonal variability accounts for up to 60% of w_e variance
240 south of the Denmark Strait, but in other regions, including both ice-covered and ice-free zones, it typically explains 10–30%.

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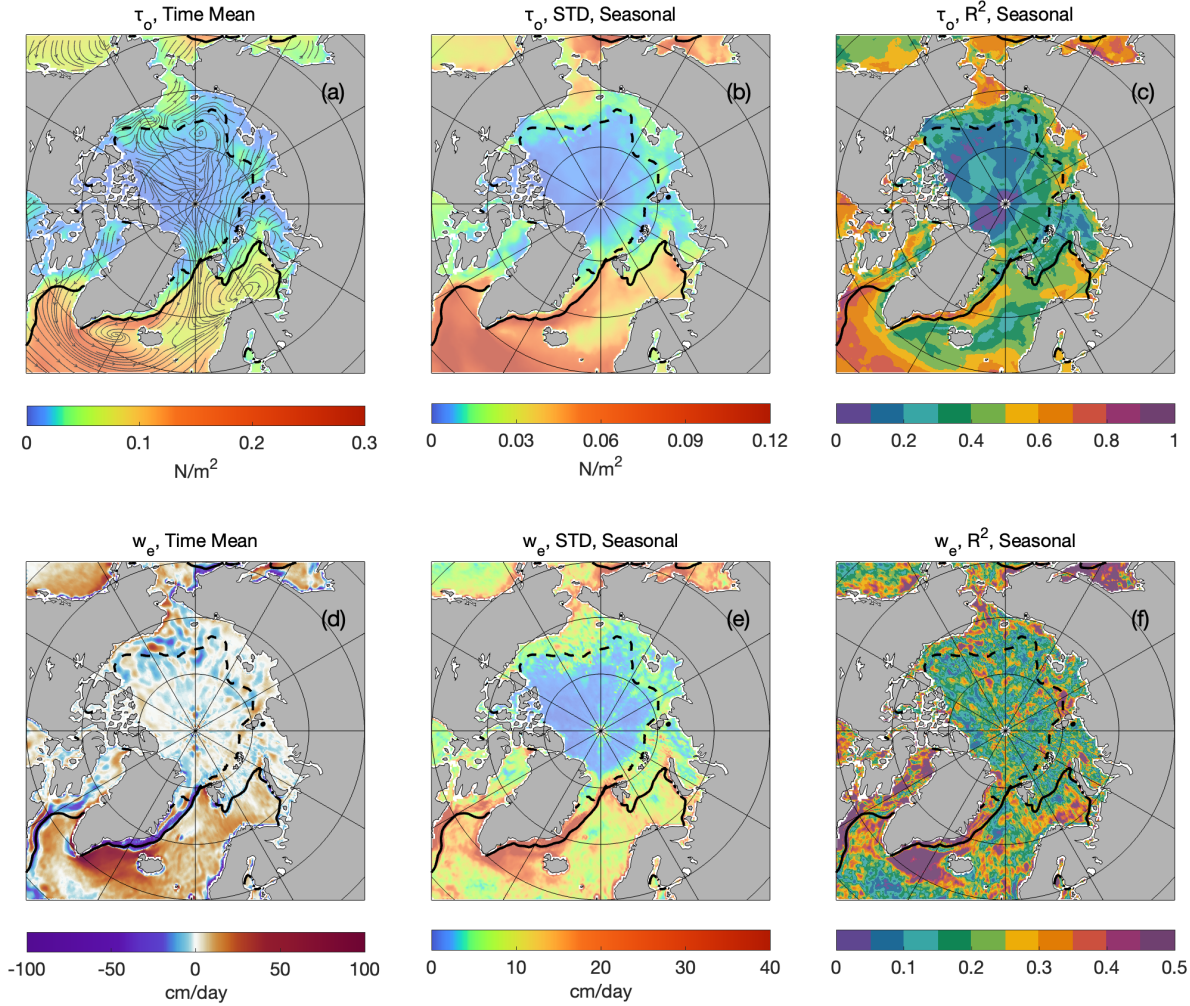


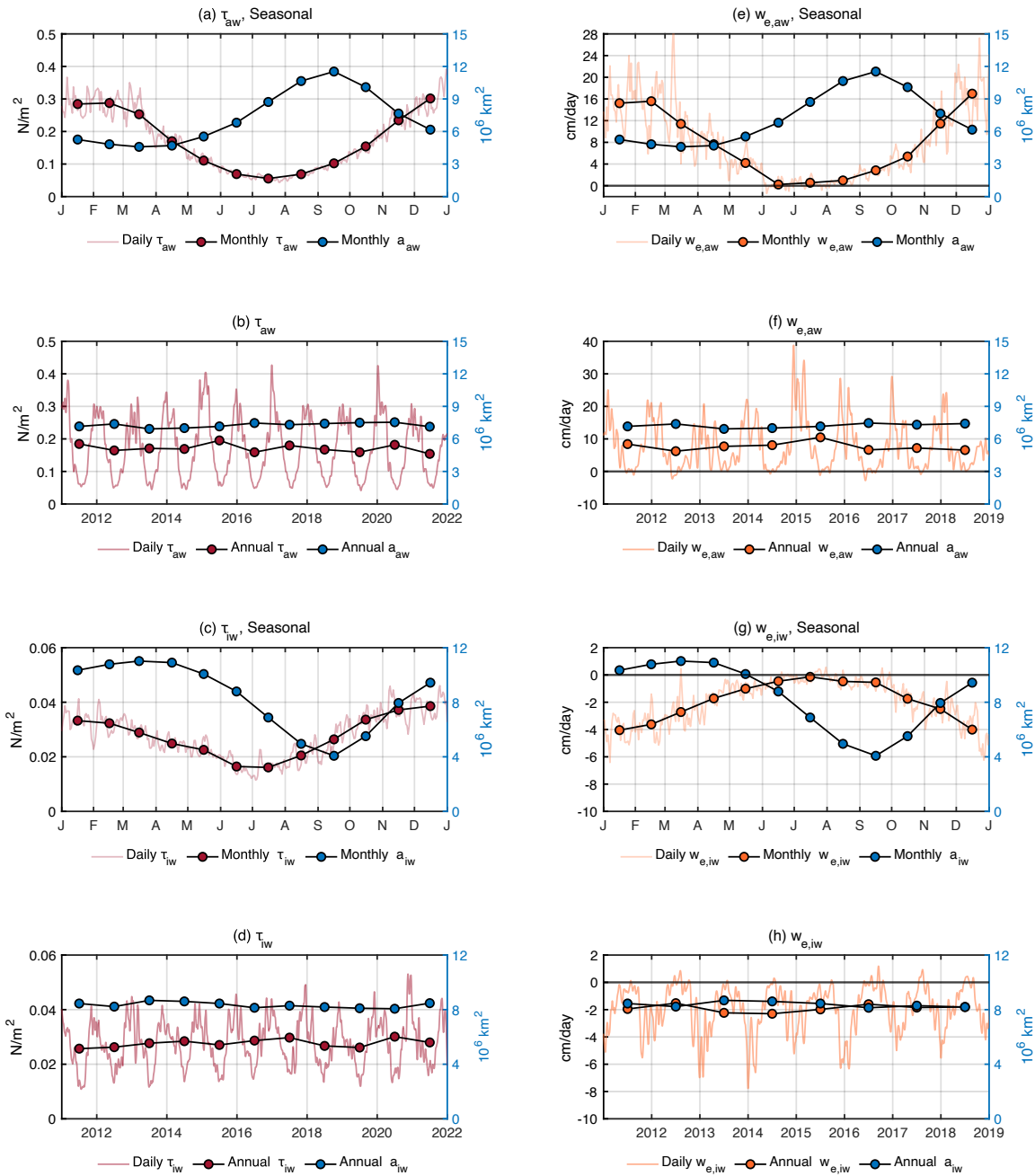
Figure 4: Mean and variability of ocean-surface stress τ_o and Ekman pumping rate w_e (positive indicates upwelling, negative indicates downwelling) in the Arctic region over 2011-2021. (a) Mean τ_o , with streamlines indicating the direction of stress. (b) Standard deviation of τ_o seasonal variability. (c) R^2 , representing τ_o variance explained by seasonal variability. (d-f) Same as (a-c) but for w_e . Streamlines in (a) show the direction of τ_o . The solid and dashed black lines represent the March and September sea ice boundaries, respectively, defined by 15% sea ice concentration averaged over 2011-2021.

The seasonal cycle of area-averaged wind-ocean surface stress (τ_{aw}) is marked by strong values in winter, peaking around 0.4 N/m^2 , and much weaker values in summer, dropping below 0.05 N/m^2 (Figure 5a). This variation corresponds to the seasonal retreat of sea ice and the associated expansion of open ocean during summer months. In ice-covered regions, the seasonal cycle of ice-ocean surface stress (τ_{iw}) is similar to that of τ_{aw} , though with significantly lower magnitudes (Figure 5c). The seasonal peak of τ_{iw} is slightly higher in 2018 than 2013, increasing from 0.010 N/m^2 to nearly 0.018 N/m^2 .

255 Due to the aforementioned uncertainties in sea ice boundary delineation, we do not present an analysis of the average Ekman
256 pumping rates w_e after 2018, as these estimates become highly sensitive to edge conditions and are thus dominated by
257 boundary artifacts.

258 In ice-free regions, the average pumping rate $w_{e,aw}$ peaks during winter upwelling, reaching around 30 cm/day, and transitions
259 to weak downwelling during the summer (Figure 5e). Annual variation in winter maximum upwelling rate is evident, with a
260 notable decline to 10 cm/day by late 2018 (Figure 5f). In contrast, in ice-covered regions, $w_{e,iw}$ is predominantly negative
261 (Figure 5g), although occasional summer upwelling events occur on daily scales. Notably, the winter downwelling rate has
262 decreased from approximately -8 cm/day in 2013–2014 to about -4 cm/day by 2017.

263



264

265 **Figure 5: Mean seasonal cycle and annual time series of area-averaged surface stress τ_o (red) and Ekman pumping rate w_e (orange,**
266 **positive indicates upwelling, negative indicates downwelling) for the Arctic region. Total areas a of the corresponding areal coverage**
267 **are also plotted in blue. Variables are subscripted aw when averaged/summed over ice-free open ocean, and are subscripted iw**
268 **when averaged/summed over ice-covered open ocean. Annual and monthly means are shown as dots in all panels. (a) Seasonal cycle**
269 **of τ_{aw} over the ice-free open ocean. (b) Timeseries of τ_{aw} from 2011-2021. (c) Seasonal cycle of τ_{iw} over the ice-covered ocean. (d)**
270 **Timeseries of τ_{iw} from 2011-2021. (e-h) same as (a-d) but for w_e and 2011-2018.**

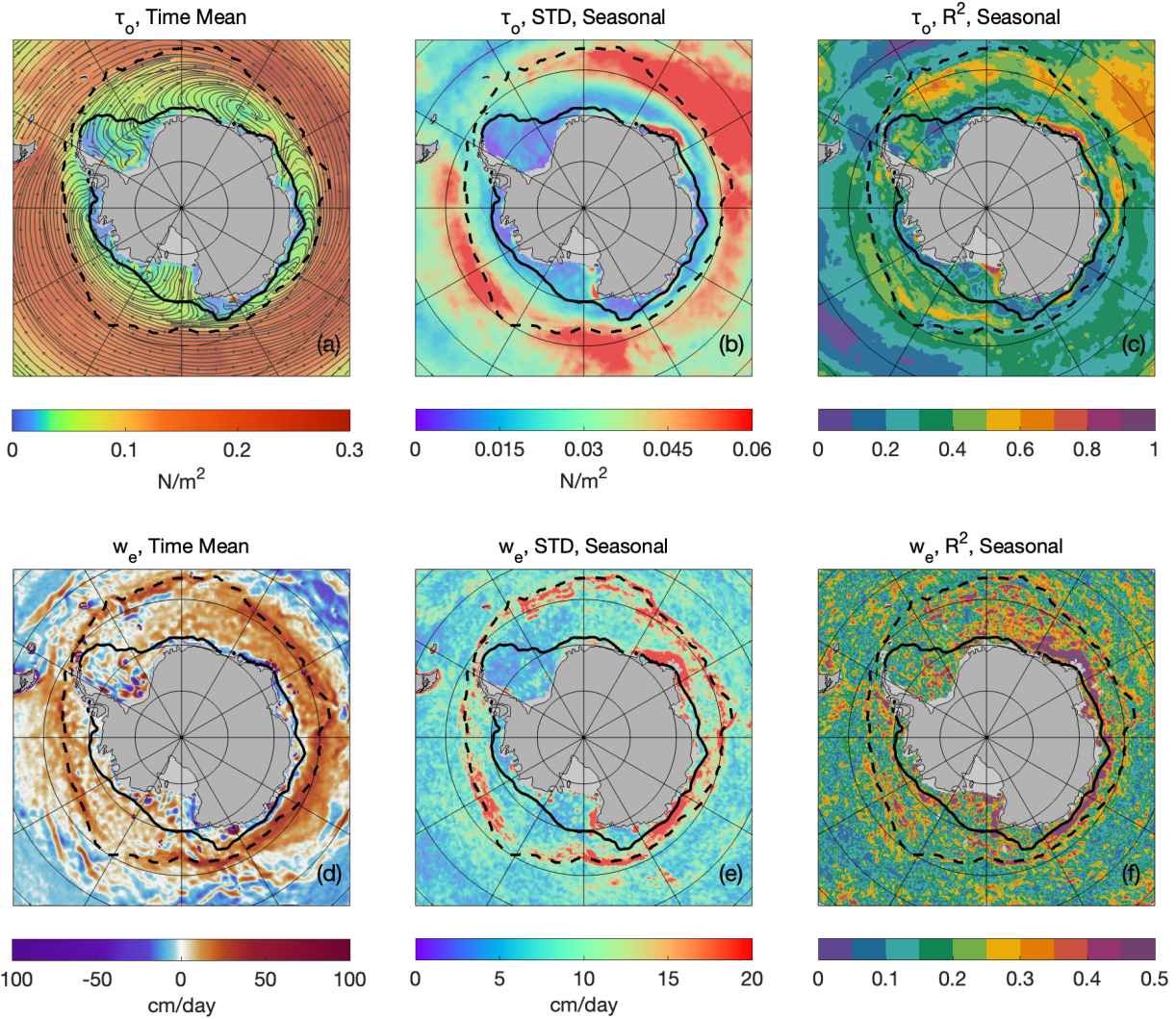
272 **3.2 Southern Ocean**

273 The spatial distribution of the time-mean τ_o in the Antarctic region is shown in Figure 6a. τ_o exhibits a prominent circumpolar
274 pattern. In ice-free regions, τ_o typically ranges from 0.2 to 0.3 N/m². In the SIZ, τ_o decreases significantly, falling to 0.04–
275 0.06 N/m², with strong regional variability.

276 The STD of τ_o seasonal variability is evidently strong near the September sea ice boundary, exceeding 0.1 N/m², particularly
277 between 0–90°E (Figure 6b). Moving northward into subpolar open-ocean, the STD gradually declines to approximately 0.04
278 N/m². Within the SIZ, seasonal variability diminishes further, typically ranging from 0.02 to 0.04 N/m². In the PIZ, it drops
279 below 0.02 N/m². The R² shows that in regions such as the Indian Ocean and southeast Pacific, seasonality explains over 50%
280 of the total variance, while in other areas, this proportion ranges from 20% to 40% (Figure 6c).

281 The spatial structure of the time-mean w_e reveals widespread upwelling south of 50°S (Figure 6d), extending nearly all the
282 way to the coast of Antarctica. In contrast to the Arctic, where strong ice-ocean coupling leads to clear transitions between
283 upwelling and downwelling across ice boundaries, the Southern Ocean does not exhibit this distinct pattern. Downwelling is
284 generally found around 55°S and farther north, or more narrowly along the Antarctic coastline.

285 The STD pattern of seasonal variability in w_e is relatively consistent across the Southern Ocean (Figure 6e), regardless of sea-
286 ice coverage, with an average value of approximately 10 cm/day. Higher variability, reaching up to 20 cm/day, occurs only
287 near the September ice boundary and is very localized. The R² pattern is also relatively homogeneous, with most areas showing
288 seasonal variability accounting for about 30% of the variance. Along the east coast of Antarctica, the seasonal cycle explains
289 more than 50% of the variance.



291

292 **Figure 6: Mean and variability of ocean-surface stress τ_o and Ekman pumping rate w_e (positive indicates upwelling, negative**
 293 **indicates downwelling) in the Antarctic region over 2013-2021. (a) Mean τ_o , with streamlines indicating the direction of stress. (b)**
 294 **Standard deviation of τ_o seasonal variability. (c) R^2 , representing τ_o variance explained by seasonal variability. (d-f) Same as (a-c)**
 295 **but for w_e . Streamlines in (a) show the direction of τ_o . The solid and dashed black lines represent the March and September sea ice**
 296 **boundaries, respectively, defined by 15% sea ice concentration averaged over 2013-2021.**

297

298 The seasonal cycle and time series of area-averaged air-water stress τ_{aw} in the Antarctic are shown in Figures 7a and 7b. In
 299 ice-free regions of the Antarctic, the average τ_{aw} peaks in August at 0.36 N/m^2 and reaches its minimum in January at 0.13
 300 N/m^2 . Annual variability is relatively small, ranging between 0.022 and 0.026 N/m^2 , with a notable positive anomaly in 2015,
 301 when the annual mean briefly increased to 0.028 N/m^2 .

302 In ice-covered regions, ice-water stress τ_{iw} shows a delayed seasonal cycle compared to τ_{aw} and peaks in September (Figure
 303 7c). It is approximately one-fifth to one-half of τ_{aw} , ranging between 0.02 N/m^2 and 0.08 N/m^2 . The seasonal pattern is

304 asymmetrical and aligns with the seasonal cycle of sea ice coverage (Eayrs et al., 2019). Similar to the Arctic, the area-averaged
305 summer minima of τ_{iw} is slightly higher in 2018 compared to 2013, increasing from 0.010 N/m² in to 0.022 N/m².
306 Before 2019, the seasonal cycle of open-ocean Ekman pumping rate $w_{e,aw}$ is relatively weak (Figure 7e), with higher values
307 in winter (12 cm/day) and lower values in summer (5 cm/day). The absence of a distinct seasonal signal is likely due to the
308 weaker seasonal cycle observed in 2017 and 2018 (Figure 7f). The annual mean varies narrowly between 7 and 9 cm/day.
309 In ice-covered regions, $w_{e,iw}$ is mostly positive throughout the year, with a brief downwelling period between January and
310 April. A shift toward stronger downwelling occurs in February, with mean values decreasing from -2 cm/day in 2013 to nearly
311 -10 cm/day by 2017. A notable anomaly occurred in 2015 when the annual mean rose sharply from 2 cm/day to 4 cm/day.
312

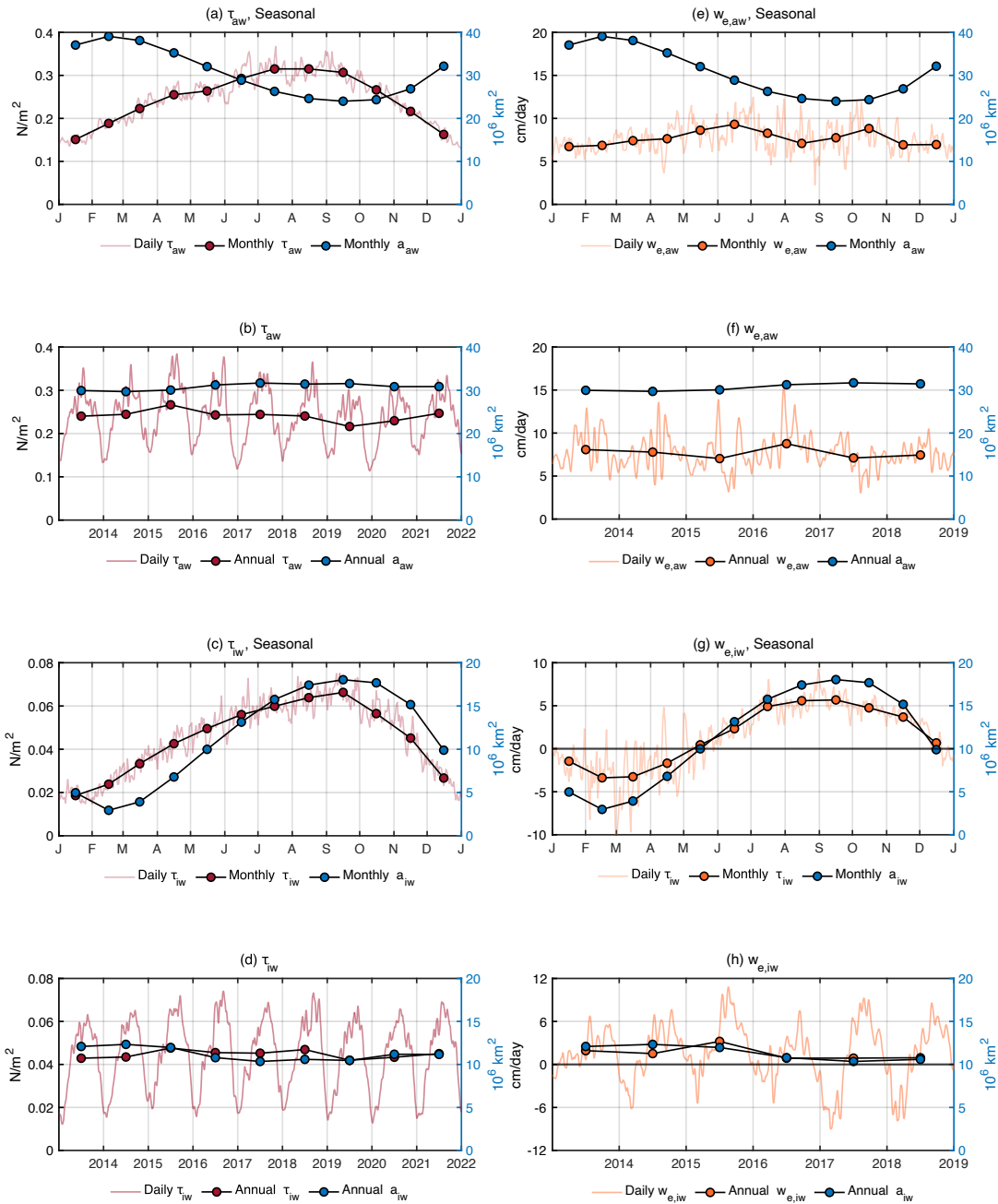


Figure 7: Timeseries and seasonal cycle of area-averaged surface stress τ_o (red) and Ekman pumping rate w_e (orange, positive indicates upwelling, negative indicates downwelling) for the Antarctic region. Total areas a of the corresponding areal coverage are also plotted in blue. Variables are subscripted aw when averaged/summed over ice-free open ocean, and are subscripted iw when averaged/summed over ice-covered open ocean. Annual and monthly means are shown as dots in all panels. (a) Seasonal cycle of τ_{aw} over the ice-free open ocean. (b) Timeseries of τ_{aw} from 2013-2021. (c) Seasonal cycle of τ_{iw} over the ice-covered ocean. (d) Timeseries of τ_{iw} from 2013-2021. (e-h) same as (a-d) but for w_e and 2013-2018.

322 4.1 Sensitivity Analysis of Ice-Water Drag Coefficient and uncertainty estimate

323 The ice-water drag coefficient, $C_{D,iw}$ is often assumed to be constant across time and space due to the scarcity of direct
 324 observations that capture its spatiotemporal variability. However, $C_{D,iw}$ can vary significantly depending on environmental
 325 conditions such as wind and wave dynamics, ice roughness, sea ice concentration, and surface morphology (Lüpkes et al.,
 326 2012; Lüpkes and Birnbaum, 2005; Cole & Stadler, 2019). Reported values for $C_{D,iw}$ range from 0.7 to over 10.0×10^{-3}
 327 (Overland, 1985; Guest and Davidson, 1987, 1991; McPhee, 2008; Cole et al., 2014), with some extreme cases reaching
 328 magnitudes on the order of 10.0×10^{-1} (Kawaguchi et al., 2024).

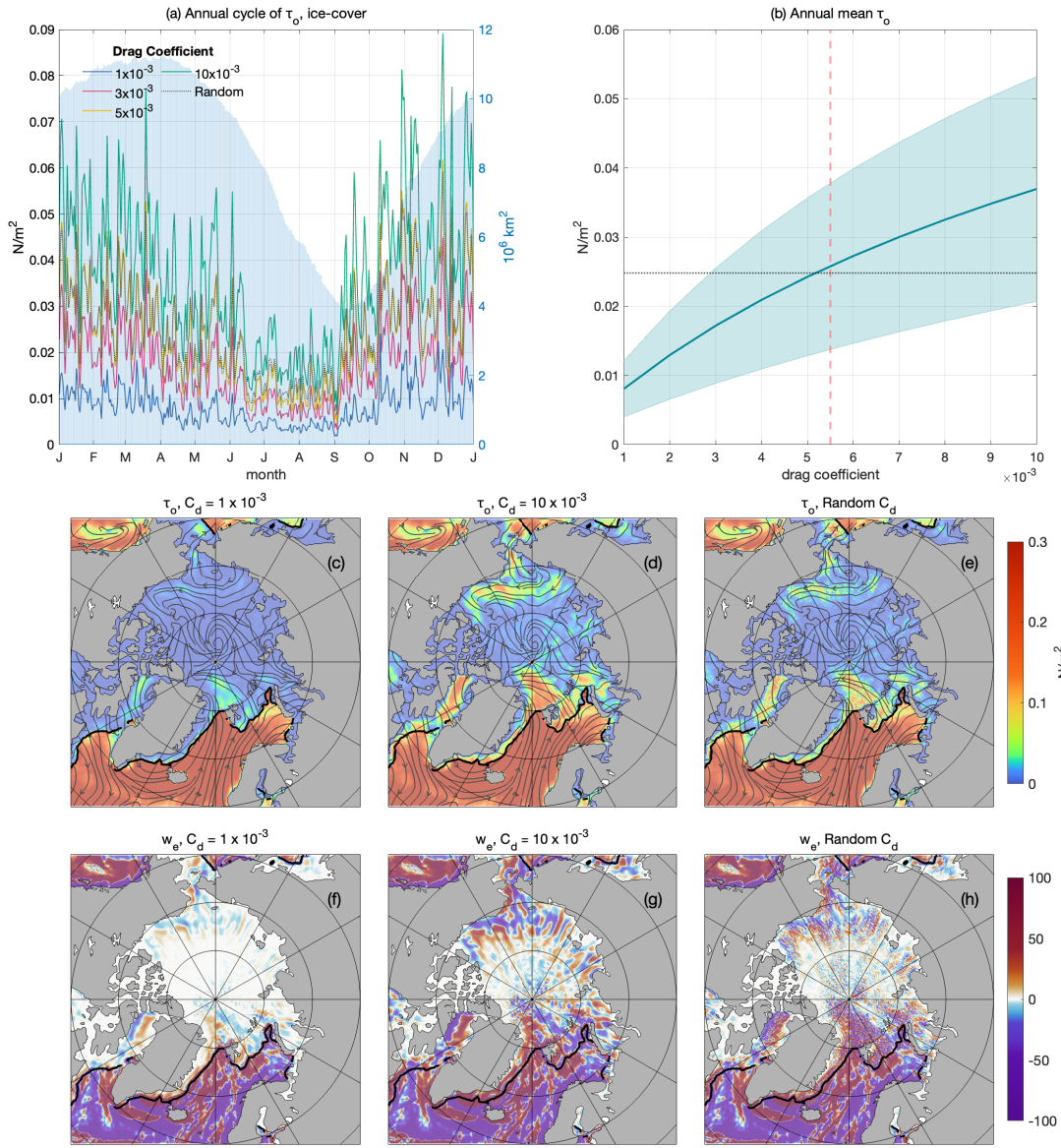
329 Commonly, a representative value of 5.5×10^{-3} has been widely adopted as a pragmatic approximation by the scientific
 330 community (Guest and Davidson, 1987; Anderson, 1987). However, this approximation may overlook important spatial and
 331 temporal variations in $C_{D,iw}$, highlighting the need for ongoing efforts to improve observations and refine its parameterization.
 332 To evaluate the sensitivity of estimated τ_o to the variations in $C_{D,iw}$, two sets of experiments are conducted for 2011: one with
 333 fixed $C_{D,iw}$ values ranging from 1.0×10^{-3} to 10.0×10^{-3} , and another using a randomized weighting map, dynamically varying
 334 $C_{D,iw}$ between between order of 10^{-3} and 10^{-2} on a daily basis at each grid cell.

335 The amplitude of τ_o scales proportionally with $C_{D,iw}$, as implied from Eq. 2 (Figure 8a). For fixed coefficients, the summer
 336 mean τ_o increases from 0.003 N/m^2 at $C_{D,iw} = 1.0 \times 10^{-3}$, to 0.015 N/m^2 at $C_{D,iw} = 10.0 \times 10^{-3}$, while winter means rise from
 337 0.012 N/m^2 to 0.053 N/m^2 . Results from the random-weighted $C_{D,iw}$ experiment closely follows the fixed cases of $C_{D,iw} = 5.0$ -
 338 6.0×10^{-3} . Similarly, the annual mean τ_o and its standard deviation increase proportionally with $C_{D,iw}$ (Figure 8b), quadrupling
 339 the annual mean and raising the standard deviation from 0.003 N/m^2 to 0.017 N/m^2 as $C_{D,iw}$ increases.

340 Figures 8c–h show the spatial distribution of τ_o and w_e in response to varying $C_{D,iw}$. Under low $C_{D,iw}$ circumstances,
 341 momentum transfer between ice and ocean is reduced, leaving small scale variability indistinct particularly in the central Arctic.
 342 As $C_{D,iw}$ increases, regions with high surface stress intensify, particularly in areas like Baffin Bay, the Chukchi Sea, and north
 343 of Fram Strait.

344 At $C_{D,iw} = 1.0 \times 10^{-3}$, the Ekman pumping rate in regions like the Fram Strait barely reach $\pm 8 \text{ cm/day}$, whereas at $C_{D,iw} = 10.0$
 345 $\times 10^{-3}$, it exceeds $\pm 30 \text{ cm/day}$, with strong contrasting upwelling and downwelling patterns. additionally, while the random-
 346 weighted $C_{D,iw}$ experiment introduces spatial noise, the broader spatial structures of both τ_o and w_e remain consistent with
 347 fixed-coefficient runs.

348



349

350 **Figure 8: Area-averaged ocean-surface stress τ_o and Ekman pumping rate w_e regarding different $C_{D,iw}$. (a) annual cycle of τ_o of**
 351 **2011, area-averaged over the sea ice-cover region. Blue areas show total ice-cover areas (right axis); (b) annual mean of τ_o with**
 352 **shading indicating one standard deviation over a year. Red dashed line marks $C_{D,iw} = 5.5 \times 10^{-3}$, black dotted line shows the annual**
 353 **mean of random $C_{D,iw}$ experiment. (c) Snapshot of τ_o with $C_{D,iw} = 1.0 \times 10^{-3}$ (3/15/2011). (d) Same as c but with $C_{D,iw} = 10.0 \times 10^{-3}$.**
 354 **(e) Same as c but with random $C_{D,iw}$. (f-h) Same as (c-e) but for w_e . Streamlines in (c-e) show the direction of τ_o . Black contours in**
 355 **(c-h) mark the 15% ice concentration on 3/15/2011.**

356

357 The final estimated uncertainty ε_{iw} , in the ice-water stress τ_{iw} is quantified daily through the integration of standard errors
 358 from sensitivity analyses of $C_{D,iw}$ and spatial Gaussian filter tests. Both filter tests and $C_{D,iw}$ tests are extended to the full

analysis period: eleven years (2011–2021) for the Arctic and nine years (2013–2021) for the Antarctic. Using the root-sum-square method, the combined uncertainty is expressed as:

$$\varepsilon_{iw} = \sqrt{(\varepsilon_{iw,F})^2 + (\varepsilon_{iw,C})^2} = \sqrt{\left(\frac{\sigma_F}{\sqrt{N_1}}\right)^2 + \left(\frac{\sigma_C}{\sqrt{N_2}}\right)^2} \quad (7)$$

where σ_F is the standard deviation of τ_{iw} from different Gaussian filter settings, and σ_C represents the standard deviation of τ_{iw} from sensitivity analysis on varying $C_{D,iw}$. The terms N_1 and N_2 denote the number of runs performed in each sensitivity analysis. This estimate assumes independence between $C_{D,iw}$ and geostrophic fields (which were spatially filtered), with perturbations of comparable amplitude between the two sets of sensitivity analysis.

Figure 9 shows the spatial distributions of relative uncertainty (ε_{iw} to τ_{iw}) for the Arctic (March 15, 2013) and the Southern Ocean (September 15, 2013) during winter season. Overall, spatial filtering produces scattered patterns (Figures 9a–9d), while varying ice-water drag coefficient yield smoother distributions (Figures 9e–9h). Median uncertainties are comparable between the two sets of experiments, ranging from 14–18% in the Arctic to 22–25% in the Antarctic. The greater uncertainties in the Antarctic reflect higher local stress variability and increased sensitivity to parameter changes, which also manifests as the higher uncertainties observed in winter compared to summer (Figures 3 and 8).

In the Arctic, combined uncertainties for zonal surface stress (τ_x) typically range from 10–20%, while locally it could exceed 100% along dynamic regions such as the Fram Strait and Beaufort Sea. Meridional stress (τ_y) exhibits similar spatial distributions, but with higher uncertainties near the Mendeleev Ridge. Median uncertainty levels for both zonal and meridional components are below 20%.

Conversely, Antarctic uncertainties are substantially higher, with median values around 40%. The highest uncertainties (>60%) are concentrated near the sea ice boundary, particularly in the eastern Weddell and Ross Seas. Regional hotspots include the Antarctic Peninsula and west of Ross Sea for τ_x , and Enderby Land and the Amundsen Sea for τ_y .

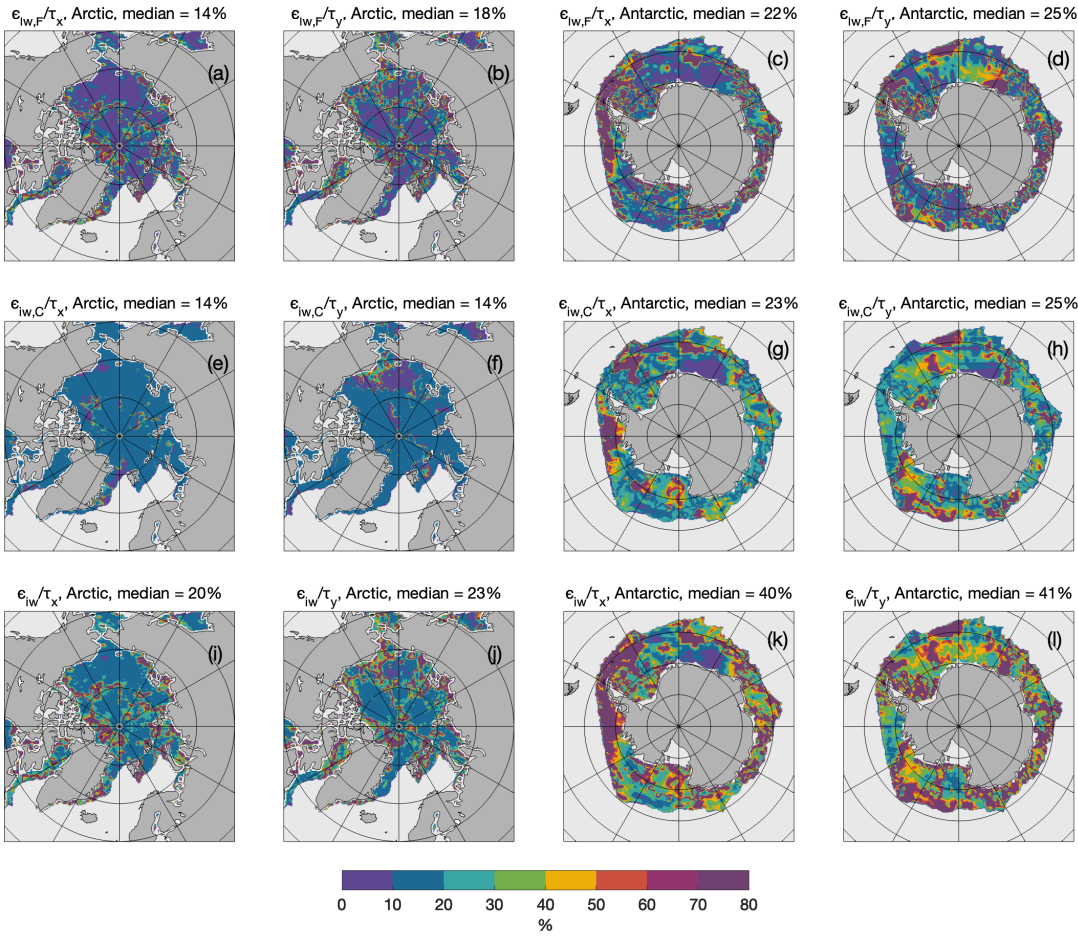


Figure 9: Estimated uncertainty fields for zonal and meridional ice-water surface stress, expressed as a ratio to the estimated ice-water surface stress in the Arctic (3/15/2013) and Southern Ocean (9/15/2013). (a) standard error introduced by Gaussian filter in the Arctic, zonal direction; (b) error from filter in the Arctic, meridional direction; (c-d) same as (a-b) but for the Antarctic; (e-h) same as (a-d) but for standard error in ice-water drag coefficient $C_{D,iw}$; (i-l) same as (a-d) but for the combined uncertainty.

In addition to the sensitivity analyses presented for drag coefficient magnitudes and spatial filtering, several other sources of uncertainty may influence the accuracy of the derived surface stress fields. First, uncertainties in the atmospheric reanalysis products used as forcing data—particularly in wind speed and direction over sea ice—can propagate directly into surface stress estimates. Prior studies have shown that different ice motion products can yield differences of at least 20–30% in polar regions due to discrepancies in boundary-layer representation, data assimilation techniques, and satellite retrieval biases (Sumata et al., 2014; Wang et al., 2022). These differences are especially pronounced in the marginal ice zone, where sharp gradients in atmospheric and surface properties are common (Wang et al., 2021; Boutin et al., 2020).

Second, the role of ocean and atmospheric stratification is not explicitly resolved in our parameterization, yet it can significantly affect stress transmission through the ice–ocean interface. Observational and modeling studies (Lüpkes & Gryanik, 2015; Lüpkes et al., 2012; Brenner et al., 2021) have shown that stability conditions in the atmospheric boundary layer modulate drag coefficients by altering turbulence and momentum fluxes—especially under stable stratification, common in winter Arctic conditions. Likewise, vertical stratification in the upper ocean can modify Ekman layer dynamics and the effective depth over which stress-induced velocities operate, introducing further uncertainty in estimates of vertical transport (Meneghello et al., 2018; Zhong et al., 2018).

Third, the spatial and temporal scales over which stresses are calculated can introduce methodological uncertainty. Coarse averaging may obscure high-frequency processes such as synoptic wind events, inertial motions, or mesoscale eddies, while finer-scale estimates risk amplifying local noise or aliasing undersampled variability (Timmermans et al., 2008; Manucharyan & Thompson, 2017; Alberello et al., 2020). This is particularly relevant in the marginal ice zone, where surface properties evolve rapidly. A more detailed analysis of scale effects and filtering sensitivity is presented in the following section.

Together, these factors point to the need for caution when interpreting surface stress magnitudes or derived quantities like Ekman pumping, particularly when used to constrain physical budgets or force ocean models. Future work should prioritize uncertainty quantification through ensemble reanalysis comparisons, the inclusion of stratification effects in drag parameterizations, and adaptive filtering techniques that respond to local dynamical conditions.

4.2 Validation with ITP Observations

Since surface stress is not usually directly measured, assessing the performance of our analysis is challenging. To address this, we revisit the assumption that surface velocity comprises both Ekman and geostrophic components, as described in Eq. (2). The geostrophic velocity (U_g) is derived from the dataset provided by Prandi et al. (2021), while the Ekman velocity component (U_e) can be easily calculated from the ocean-surface stress using Eq. (4).

This assumption provides a first-order approximation of surface velocity, and neglects other processes, such as ageostrophic motions, vertical shear, and submesoscale dynamics, which may introduce additional uncertainties. To robustly validate satellite-derived ocean surface stress estimates, we compared the derived surface velocity, i.e., sum of U_g and U_e , with in situ measurements from Ice-Tethered Profilers (ITPs, Krishfield et al., 2008; Toole et al., 2011; <http://www.whoi.edu/itp>). In particular, several ITPs that equipped velocity sensors (ITP-V, Williams et al., 2010) are used. To align the temporal resolution of the datasets, we processed the ITP data by computing daily and weekly means, facilitating direct comparisons with daily satellite products. This approach avoids the uncertainties associated with interpolating satellite data to match the higher-frequency ITP profiles, which could introduce significant errors due to the under-sampling nature of satellite observations.

Despite this temporal alignment, inherent limitations persist due to spatial and temporal sampling discrepancies. ITPs provide high-resolution vertical profiles at specific locations, capturing fine-scale and transient oceanic features. In contrast, satellite observations offer broader spatial coverage but may not resolve such fine-scale variability, especially in polar regions where

data gaps are common due to persistent cloud cover and sea ice. These differences can lead to reduced correlation and increased bias in validation statistics, as observed in previous studies comparing satellite-derived sea surface salinity products with in situ observations (Thouvenin-Masson et al., 2022; Boutin et al., 2016; Vinogradova & Ponte, 2013). It is important to note that this comparison does not serve as a definitive validation of the absolute accuracy of our stress estimates. Instead, it assesses whether the foundational assumptions underpinning our analysis sufficiently represent the complex dynamics of the Arctic Ocean.

We use velocity data collected from five ITP-V missions deployed on multiyear sea ice in the Canada Basin between 2011 and 2019 (Figure 10; Table 2). Observations from ITP-77, ITP-78, and ITP-79 are truncated to exclude periods with significant data gaps and drifts near the end of their deployments.

The five ITP-Vs are categorized into two groups based on deployment timing and drift trajectories. ITP-70 and ITP-80 were deployed during summer, operated for ~300 days, and primarily drifted between 75–80°N. In contrast, ITP-77, ITP-78, and ITP-79, deployed in March 2014, operated for less than 200 days and followed more constrained east-to-west trajectories between 73–75°N.

439

440 **Table 2: Details of the ITP-V records**

Unit ID	Start		Last		# of Days	# of Profiles
	Position	Date	Position	Date		
ITP-70	76.81°N	8/26/2013	77.11°N	7/15/2014	324	3713
	138.89°W		156.51°W			
ITP-77	73.37°N	3/11/2014	75.89°N	10/2/2014	206	2367
	134.99°W		158.50°W		(153*)	(1800*)
ITP-78	74.36°N	3/12/2014	74.08°N	8/6/2014	148	1694
	135.14°W		145.43°W		(130*)	(1500*)
ITP-79	75.38°N	3/22/2014	75.02°N	9/30/2014	193	1694
	136.50°W		148.37°W		(143*)	(1636*)
ITP-80	77.36°N	8/14/2014	75.68°N	5/24/2015	284	3260
	146.15°W		151.79°W			

441 *Data towards the end of the series exhibits quality issues that necessitate truncation.

442

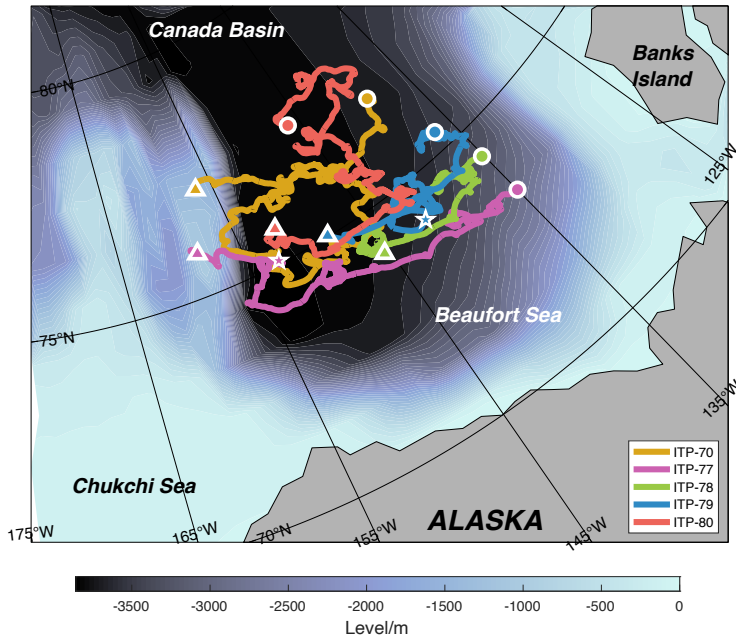


Figure 10: ITP-V drift paths in the Arctic Ocean (colored curves). The deployment locations are marked by circles, latest locations by triangles, and the cutoff locations for ITP-77 and ITP-79 by stars.

To account for temporal sampling differences and mitigate aliasing from high-frequency variability, the sub-daily ITP-V velocity data are first averaged to daily means before comparison with the satellite-derived velocity field ($U_g + U_e$). The corresponding satellite values are then extracted at the nearest grid point along each ITP track (Figure 11). This approach reduces mismatch due to subsampling in the satellite product and ensures a more consistent temporal basis for comparison. Along the path of ITP-70 and ITP-80, satellite-derived ocean surface velocities exhibit moderate agreement with in situ observations, particularly in capturing high-frequency variability. In contrast, comparisons with ITP-77, ITP-78, and ITP-79 reveal weaker correspondence, most notably in the zonal velocity components. For the meridional component, ITP-77 shows relatively better alignment during the initial ~100-day period until mid-July.

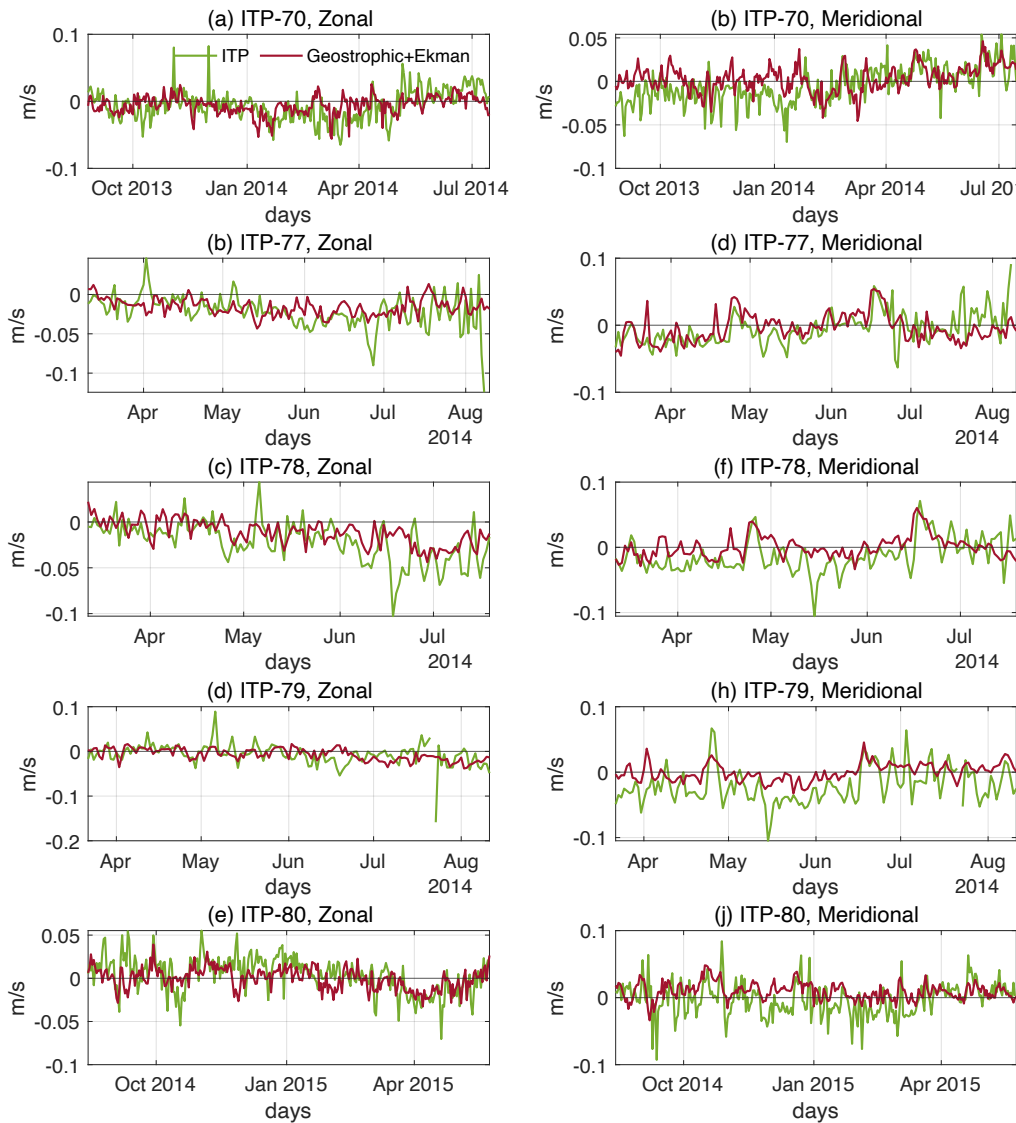


Figure 11: Daily mean timeseries of zonal and meridional surface velocity. (a) Zonal velocity along ITP-70 paths. (b) Zonal velocity along ITP-77 paths. (c) Zonal velocity along ITP-78 paths. (d) Zonal velocity along ITP-79 paths. (e) Zonal velocity along ITP-80 paths. (f-j) same as (a-e) but for meridional velocity. Green curves represent velocity data retrieved from ITPs. Red curves are collocated obtained from the satellite-derived velocity fields, i.e., geostrophic plus Ekman velocity.

Figure 12 presents the comparison of satellite-derived surface velocity components against collocated ITP-V observations for the ITP-70 (panels a,b) and ITP-80 (panels c,d) paths. For ITP-70, the zonal component yields a Pearson correlation of $r = 0.31$ and standard deviation of 0.022 m s^{-1} , while the meridional component gives $r = 0.42$ and $\text{STD} = 0.020 \text{ m s}^{-1}$. ITP-80 exhibits slightly stronger zonal agreement ($r = 0.43$) but weaker meridional agreement ($r = 0.34$). In both deployments, scatter

markedly decreases for observations taken after ~200 days (warm colors), accompanied with the predominantly northward drift of ITP-70 and westward drift of ITP-80 seen in Figure 10.

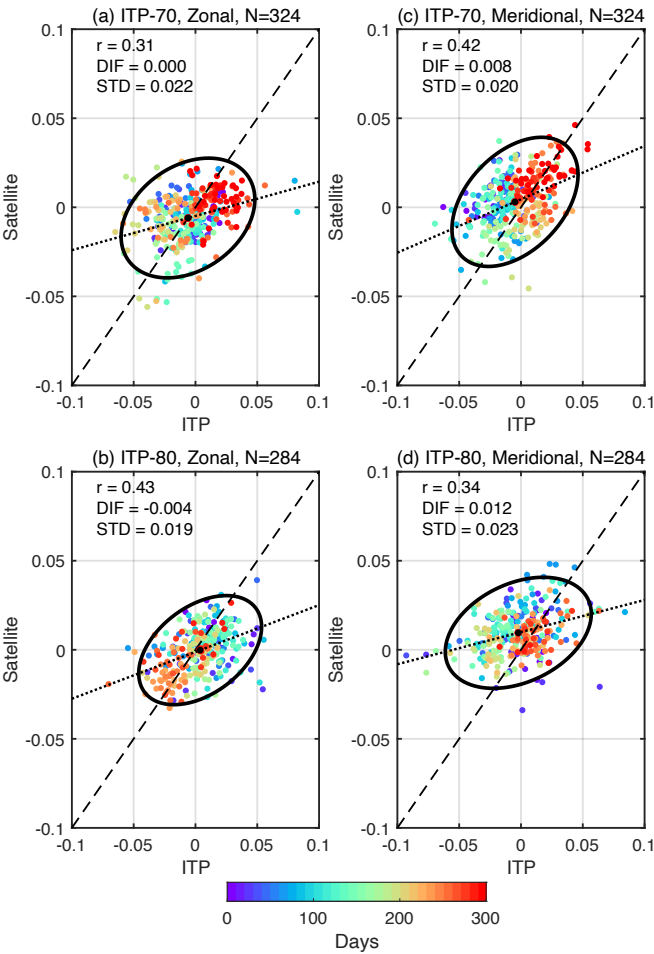


Figure 12: Scatterplots of collocated surface velocity pairs for ITPs with data spanning more than 200 days (unit: m/s). (a) Zonal velocity along ITP-70 paths. (b) Zonal velocity along ITP-80 paths. (c-d) Same as (a-b) but for meridional velocity. The total number of days (N) is given. Correlation coefficients, mean differences (DIF), and standard deviations (STD) of the differences between satellite-derived velocity and ITP observations are also displayed. 95% confidence ellipse (black contour), linear fitting (black dotted line) are also given in each panel.

By contrast, Figure 13 summarizes the additional ITP observational periods (panels a–f), where correlation coefficients span $r = 0.07\text{--}0.56$, and $\text{STD} = 0.021\text{--}0.025\text{ m s}^{-1}$. One component reaches a moderate correlation ($r = 0.56$), while most remain weak ($r < 0.40$), and no coherent temporal clustering is apparent. Across all deployments, satellite-derived velocities exhibit a slight northward bias and tend to underestimate ITP-measured surface speeds beyond 100 days post-release.

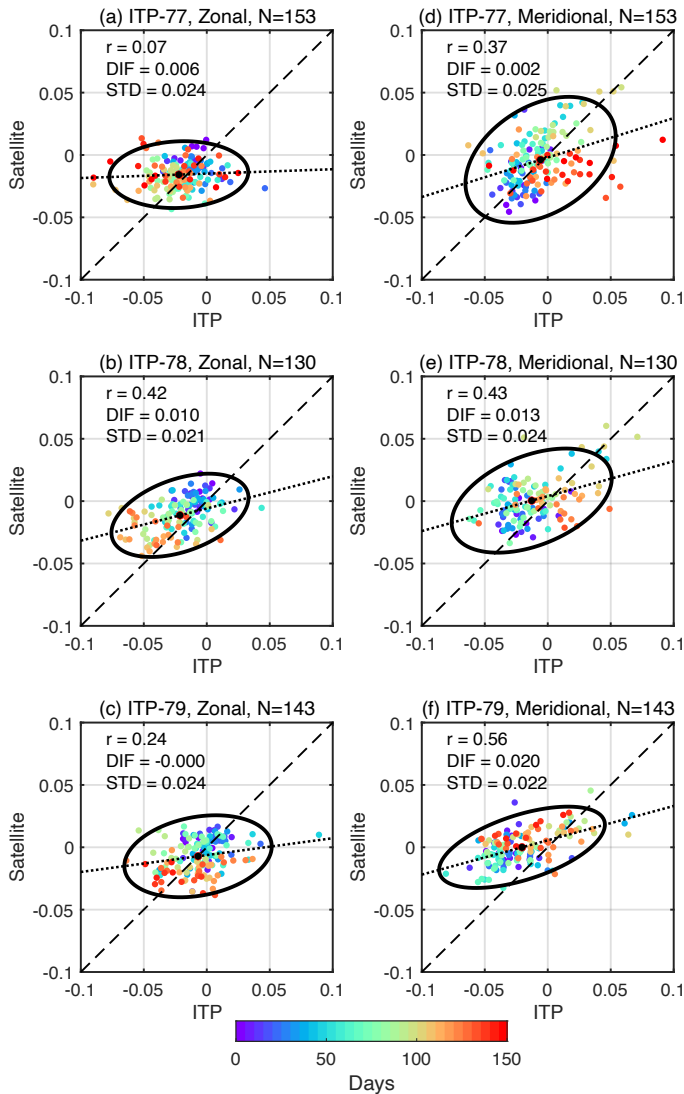


Figure 13: Scatterplots of collocated surface velocity pairs for ITPs with data spanning less than 200 days (unit: m/s). (a) Zonal velocity along ITP-77 paths. (b) Zonal velocity along ITP-78 paths. (c) Zonal velocity along ITP-79 paths. (d-f) Same as (a-c) but for meridional velocity. The total number of days (N) is given. Correlation coefficients, mean differences (DIF), and standard deviations (STD) of the differences between satellite-derived velocity and ITP observations are displayed. 95% confidence ellipse (black contour), linear fitting (black dotted line) are also given in each panel.

Table 2 and Figure 14 provide a comprehensive comparison between satellite-derived surface velocity estimates and in situ velocity measurements obtained from ITP-V across five deployments. The analysis includes both the zonal (east–west) and meridional (north–south) velocity components and considers statistics derived from both daily and weekly averaged timeseries. A consistent northward bias is evident in the satellite-derived velocities across most ITP paths. While the mean zonal velocities are generally in close agreement between satellite and ITP-V data—for instance, at ITP-70 both sources report a mean of –

493 0.006 m/s—some deployments exhibit a more pronounced bias. ITP-77 and ITP-78, in particular, show a noticeable eastward
 494 offset, with satellite-derived zonal velocities being more positive than the ITP-V counterparts by approximately 0.005–0.01
 495 m/s. The meridional component aligns more closely overall, but satellite velocities still tend to be more northward. These
 496 biases are evident in the top-left panel of Figure 14, where most points lie above the 1:1 line, especially in the meridional
 497 direction.

498 In addition to the bias, satellite-derived velocities systematically exhibit reduced variability compared to ITP-V observations.
 499 Across all deployments and components, the STD of satellite velocity is consistently lower than that observed in the ITP-V
 500 data. For example, while the average zonal STD in the ITP data is around 0.022 m/s, the corresponding satellite value is
 501 approximately 0.011 m/s. The top-right panel of Figure 14 illustrates this discrepancy clearly: all data points fall below the 1:1
 502 line, indicating that satellite products underestimate temporal variability. This reduced variability likely reflects the filtering
 503 and smoothing inherent in satellite altimetry products, which are designed to represent large-scale geostrophic flows and may
 504 not fully resolve the higher-frequency or smaller-scale fluctuations captured by the ITP instruments.

505 Despite these limitations, satellite-derived velocity fields are able to explain a substantial portion of the observed variance in
 506 the ITP-V measurements on daily mean scale. The coefficient of determination (R^2) values indicate that, on average, satellite
 507 products account for about 50% of the variability in the daily data. This explanatory power increases substantially when the
 508 analysis is performed on weekly averaged time series, with R^2 reaching as high as 0.77 for the zonal component at ITP-78 and
 509 0.60 for the meridional component at ITP-79. These results suggest that, although satellite estimates smooth out finer-scale
 510 variability, they capture the dominant patterns in large-scale motion effectively. However, the correlation coefficients between
 511 satellite and ITP-V velocities are more modest, typically ranging from 0.3 to 0.4 in the daily records. With weekly averaging,
 512 these correlations improve significantly, sometimes by more than 0.2, as illustrated in the bottom-left panel of Figure 14. This
 513 panel shows weekly-averaged points clustering nearer to or above the diagonal, particularly for ITP-78 and ITP-79, while daily
 514 correlations tend to remain lower and more scattered.

515 The performance of the satellite velocity estimates varies by deployment. ITP-78 and ITP-79 demonstrate the strongest
 516 agreement. For example, ITP-78’s zonal component yields correlation coefficients of 0.42 for daily data and 0.78 for weekly
 517 data, with R^2 values of 0.32 and 0.77, respectively. Similarly, ITP-79’s meridional component shows a daily correlation of
 518 0.56 and a weekly correlation of 0.76, with corresponding R^2 values of 0.47 and 0.60. These high values underscore the ability
 519 of satellite altimetry to capture meaningful geophysical signals under favorable conditions. Conversely, performance is notably
 520 weaker at ITP-77, where the zonal velocity component yields a daily correlation of only 0.07 and a weekly correlation of 0.31,
 521 suggesting a diminished ability of satellite products to resolve local variability in that particular region or time frame. Such
 522 differences likely arise from a combination of regional oceanographic complexity and satellite data limitations, including issues
 523 related to proximity to sea ice or the presence of submesoscale activity not well resolved by gridded products.

524

525 **Table 2: Comparison of daily satellite-derived velocity to the ITP velocity along the ITP tracks. Correlations (r) with $p < 0.05$ are in**
 526 **bold. Numbers in the bracket are from the weekly mean timeseries.**

ITP-Vs		Mean (ITP)	Mean (Sat.)	STD (ITP)	STD (Sat.)	<i>r</i>	R ²
ITP-70	Zonal	-0.006	-0.006	0.022	0.014	0.31 (0.45)	0.12 (0.32)
	Meridional	-0.005	-0.003	0.021	0.015	0.42 (0.59)	0.20 (0.40)
ITP-77	Zonal	-0.022	-0.016	0.023	0.011	0.07 (0.31)	0.11 (0.20)
	Meridional	-0.005	-0.004	0.024	0.021	0.37 (0.47)	0.01 (0.02)
ITP-78	Zonal	-0.021	-0.011	0.022	0.014	0.42 (0.78)	0.32 (0.77)
	Meridional	-0.013	0.001	0.026	0.017	0.43 (0.61)	0.25 (0.32)
ITP-79	Zonal	-0.006	-0.007	0.024	0.013	0.24 (0.38)	0.02 (0.11)
	Meridional	-0.020	-0.001	0.027	0.013	0.56 (0.76)	0.47 (0.60)
ITP-80	Zonal	-0.004	-0.001	0.020	0.012	0.42 (0.61)	0.30 (0.60)
	Meridional	-0.002	0.010	0.024	0.012	0.34 (0.52)	0.23 (0.16)
Mean	Zonal	-0.010	-0.008	0.022	0.011	0.29 (0.50)	0.13 (0.40)
	Meridional	-0.009	0.002	0.023	0.016	0.42 (0.59)	0.22 (0.28)

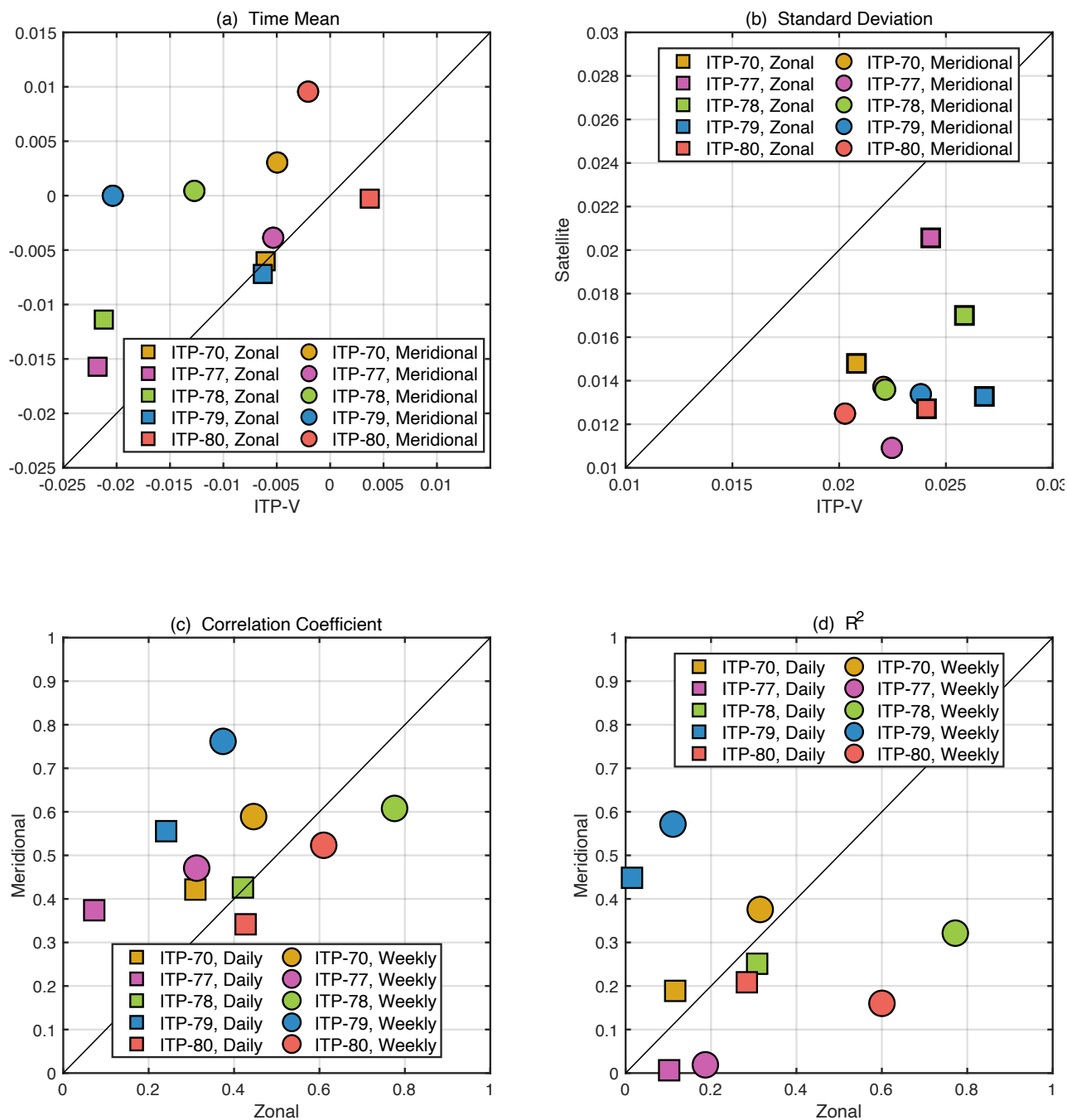


Figure 14: Scatterplots of statistics of satellite-derived velocity and the ITP velocity. (a) Temporal mean and (b) standard deviation of the paired velocity (unit: m/s). (c) Correlation between satellite-derived velocity and the ITP velocity. (d) Coefficient of determination (R^2) of the variation in the ITPs explained by satellite-derived velocity.

533 The comparison between satellite-derived ocean surface stress estimates and ITP observations, despite their differences in
534 resolution and sampling, shows encouraging agreement at the first-order level. The satellite stress products successfully capture
535 the broad spatial and temporal variability in surface velocity, supporting the utility of satellite-based estimates in reflecting
536 first-order dynamical signals. This level of agreement supports the overall utility of satellite products in characterizing large-
537 scale stress variability and motivates their continued use in data-sparse polar regions.

538 However, key limitations remain due to inherent mismatches in spatial and temporal sampling between the datasets. Satellite
539 observations, with their typical resolution of ~25 km and daily sampling frequency, cannot resolve the submesoscale variability
540 and high-frequency processes detectable in the pointwise and often sub-daily ITP profiles (Timmermans et al., 2008). This
541 spatial averaging can smooth out gradients in wind, ice motion, or stress fields that may be sharply defined at smaller scales,
542 particularly in regions such as the marginal ice zone (MIZ) where sea ice concentration and morphology are highly variable
543 (Manucharyan & Thompson, 2017; Alberello et al., 2020).

544 Temporally, satellite-derived surface stress products may fail to capture transient forcing events such as storm-driven
545 accelerations, inertial oscillations, or short-lived leads in sea ice. In contrast, ITPs can resolve such high-frequency processes
546 (Toole et al., 2011; Timmermans et al., 2012), leading to potential discrepancies when aligning the two datasets. Furthermore,
547 since satellite estimates of stress are often derived from independent wind and ice motion products (e.g., National Centers for
548 Environmental Prediction Reanalysis, NSIDC drift), their accuracy is subject to the limitations of those input fields (Sumata
549 et al., 2014; Lavergne et al., 2010). The lack of fully collocated wind and ice motion fields at the exact time and location of
550 ITP measurements compounds the uncertainty.

551 These spatial and temporal mismatches further introduce representation errors, as mismatches are not due to sensor or
552 algorithm flaws but rather due to sampling disparities (Janjić et al., 2008). Such errors have been well documented in the
553 context of satellite sea surface salinity (Boutin et al., 2016; Vinogradova et al., 2019). These limitations are evident in the
554 reduced agreement observed for ITP-77, ITP-78, and ITP-79, where several compounding factors likely contributed. First, the
555 ~25 km resolution of the satellite product may be insufficient to resolve submesoscale features and sharp velocity gradients.
556 Second, the timing of deployment in March overlaps with a period of elevated kinetic energy in the Beaufort Gyre (Cassianides
557 et al., 2023), during which intensified eddy activity in the Canada Basin enhances mesoscale variability (Son et al., 2022;
558 Regan et al., 2020). This variability, well captured by the high-resolution ITP profiles, is easily aliased or smoothed out in the
559 satellite-derived daily fields, further amplifying mismatches in direct comparisons.

560 Moving forward, best practices in validation should account for these differences explicitly. The development of higher-
561 resolution satellite products (Auger et al., 2022; Lucas et al., 2023), along with assimilation into coupled models (Wang et al.,
562 2018), also offers a promising path forward. Increased density of ITP deployments, moored arrays, and coordinated airborne
563 campaigns (Perovich et al., 2023) will be crucial for better spatial coverage in dynamic regions like the Beaufort Gyre and the
564 MIZ.

565 In conclusion, the comparison reveals that while satellite-derived velocities are subject to systematic biases and reduced
566 variability relative to in situ observations, they nonetheless capture a significant portion of the observed variance, particularly

when considered at weekly timescales. The agreement is stronger for the meridional component and in regions where large-scale geostrophic flows dominate. These results support the use of satellite-derived velocity products for basin-scale circulation studies, while also highlighting the need for caution in applications requiring high-frequency or fine-scale flow resolution.

5 Data Availability

Daily fields of ocean-surface stress vectors and derived vertical Ekman velocity for the polar oceans are provided for two periods: 2011–2021 for the Arctic (EPSG number 3408) and 2013–2021 for the Antarctic (EPSG number 3409) and are available at <https://doi.org/10.5281/zenodo.15534576> (Liu & Yu, 2024). The datasets include three auxiliary fields: (i) land mask, (ii) grid longitudes and latitudes, and (iii) uncertainty estimates for ocean-surface stress.

The input datasets can be found at NSIDC (ice motion: <https://nsidc.org/data/nsidc-0116/versions/4>; ice extent: <https://nsidc.org/data/nsidc-0051/versions/2>) and AVISO website (dynamic topography: <https://www.aviso.altimetry.fr/en/data/products/sea-surface-height-products/regional/arctic-ocean-sea-level-heights.html>). ITP-V data used in this work are retrieved from the WHOI website at <https://www2.whoi.edu/site/itp/>. CPOM-DOT/geostrophic currents data are provided by the Centre for Polar Observation and Modelling, University College London (https://www.cpom.ucl.ac.uk/dynamic_topography). The associated scripts and packages used in this study are openly available on GitHub at (https://github.com/cydenyliu/Polar_Stress).

6 Conclusions

This work presents a daily, 25 km resolution dataset of satellite-derived ocean-surface stress for the Arctic (2011–2021) and Southern Oceans (2013–2021). The dataset provides detailed daily maps of τ_o across polar regions north of 60°N and south of 50°S. This dataset achieves finer spatial and temporal resolution, enabling more precise analysis of short-term air-sea interactions and regional Ekman dynamics. In both the Arctic and Antarctic, it captures short-term and sharp transition between Ekman upwelling in ice-free regions and downwelling in ice-covered areas.

Uncertainty in the derived ocean-surface stress fields arises primarily from two sources. The first is the spatial filter applied to the SSH datasets, which reduces small scale variability and enhances consistency between the sea level fields. The second source of uncertainty is related to the ice-water drag coefficient, which is poorly observed and can vary significantly between order of 10^{-3} and 10^{-2} . These factors result in a median uncertainty of approximately 20% in the Arctic and about 40% in the Southern Ocean.

The derived Ekman velocity is used to validate against ITP data from the Arctic's Canada Basin. Satellite-derived surface velocity, which combine Ekman and geostrophic components, capture over 50% of the observed variation in surface velocity. Correlation coefficients range from 0.6 to 0.8 on monthly and longer timescales, indicating moderate to strong agreement. It is important to consider the complex dynamics of the Arctic Ocean when interpreting these statistics. In addition to Ekman

597 and geostrophic velocity (Regan et al., 2019), processes such as shallow eddy activity (Timmermans et al., 2008; Kenigson et
598 al., 2021; Meneghello et al., 2021), turbulent mixing (Guthrie et al., 2013; Kawaguchi et al., 2014, 2019), and internal waves
599 (Kawaguchi et al., 2016; Zhao et al., 2016) also contribute to the observed variability. Many of these processes remain
600 challenging to observe and parameterize.

601 While atmospheric reanalysis products such as ERA5 offer wind estimate in the polar regions, they do not explicitly capture
602 sea ice interaction, and coupled models, though detailed, are computationally expensive and often opaque in their assumptions.
603 Our product bridges this gap by offering a reproducible, observationally constrained dataset that supports process studies and
604 model validation. Despite some simplifying assumptions, it has comparably spatial resolution to ERA5 over open ocean and
605 offers added value in sea ice regions.

606 Future updates will focus on two primary areas. We plan to extend the dataset’s temporal coverage through 2021 by
607 incorporating updated versions of OAFlux and other relevant data products as they become available. This will ensure
608 consistency across components while maintaining the dataset’s reliability. Second, the availability of reliable surface height
609 products for the polar region will further enhance data accuracy. While awaiting these advancements, we will assess the
610 potential impacts of transitioning to reanalysis data on our results. Additionally, future research will address key processes that
611 remain underrepresented, such as variable Ekman depth and mesoscale turbulence, to refine the depiction of polar ocean
612 dynamics. Incorporating these factors will improve the ability to capture localized features critical for understanding air-ice-
613 ocean interactions.

615 **Appendix A: Glossary of Terminology**

616 **Table A1: glossary of terminology and acronyms used in this study.**

<i>Terminology/Acronyms</i>	<i>Description</i>
EASE	Equal-Area Scalable Earth grid
SSH	Sea Surface Height
OAFlux	Objectively Analyzed Air-Sea Fluxes
NSIDC	National Snow and Ice Data Center
MEaSUREs	Making Earth System Data Records for Use in Research Environments
COARE	Coupled Ocean–Atmosphere Response Experiment
SSM/I	Special Sensor Microwave/Imager
AMSR-E	Advanced Microwave Scanning Radiometer for EOS
AVHRR	Advanced Very High Resolution Radiometer
MODIS	Moderate Resolution Imaging Spectroradiometer

QuikSCAT	Quick Scatterometer
IABP	International Arctic Buoy Programme
CLS/PML	Collecte Localisation Satellites / Plymouth Marine Laboratory
SMMR	Scanning Multichannel Microwave Radiometer
DMSP	Defense Meteorological Satellite Program
SSM	Special Sensor Microwave
I-SSMIS	Improved Special Sensor Microwave Imager/Sounder
RMSD	Root Mean Square Deviation
CPOM	Centre for Polar Observation and Modelling
DOT	Dynamic Ocean Topography
SIZ	Seasonal Ice Zone
PIZ	Perennial Ice Zone
MIZ	Marginal Ice Zone
STD	Standard Deviation
R ²	Coefficient of Determination
ITP	Ice-Tethered Profiler

617

618 **Author contributions**

619 CL: conceptualization, data curation, formal analysis, methodology, software, visualization, writing – original draft preparation,
620 writing – review and editing. LY: conceptualization, project administration, supervision, validation, writing – review and
621 editing.

622 **Competing interests**

623 The contact author has declared that none of the authors has any competing interests.

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