

A synthetic [optical](#) database generated by radiative transfer simulations in support of studies in ocean optics and optical remote sensing of the global ocean

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Abstract. Radiative transfer (RT) simulations have long been used to study relationships between the inherent optical properties (IOPs) of seawater and light fields within and leaving the ocean from which the ocean apparent optical properties (AOPs) can be calculated. For example, inverse models to estimate IOPs from ocean color radiometric measurements have been developed or validated using results of RT simulations. Here we describe the development of a new synthetic optical database based on hyperspectral RT simulations across the spectral range from the near-ultraviolet to near-infrared performed with the HydroLight radiative transfer code. The key component of this development was the generation of the synthetic dataset of seawater IOPs ~~which that~~ served as input to RT simulations. Compared to similar developments of optical databases in the past, the present dataset of IOPs is characterized by probability distributions of IOPs that are consistent with global distributions representative of vast areas of open ocean pelagic environments and coastal regions covering a broad range of optical water types. The generation of the synthetic data of IOPs associated with particulate and dissolved constituents of seawater was driven largely by an extensive set of field measurements of the phytoplankton absorption coefficient collected in diverse oceanic environments. Overall, the synthetic IOP dataset consists of 3320 combinations of IOPs. Additionally, the pure seawater IOPs were assumed following recent recommendations. The RT simulations were performed using 3320 combinations of input IOPs assuming vertical homogeneity within an infinitely deep ocean. These input IOPs were used in three simulation scenarios associated with assumptions about inelastic radiative processes in the water column ([not considered in previous synthetically-generated optical databases](#)) and three simulation scenarios associated with sun zenith angle. Specifically, the simulations were made assuming no inelastic processes, the presence of Raman scattering by water molecules, and the presence of both Raman scattering and fluorescence of chlorophyll-a pigment. Fluorescence of colored dissolved organic matter was omitted from all simulations. For each of these three simulation scenarios, the simulations were made for three sun zenith angles of 0°, 30, and 60° assuming clear skies, standard atmosphere, and wind speed of 5 m s⁻¹. Thus, overall 29880 RT simulations were performed. The output results of these simulations include the radiance distributions, plane and scalar irradiances, and the whole set of AOPs including the remote-sensing reflectance, vertical diffuse attenuation coefficients, and mean cosines where all optical variables are reported in the spectral range from 350 to 750 nm at 5 nm intervals for different depths between the sea surface and 50 m. The consistency of this new synthetic database has been assessed through comparisons with in situ data and previously developed empirical relationships involving the IOPs and AOPs. [The database is available at Dryad open-access repository of research data \(doi:10.6076/D1630T\).](#)

41

42 1 Introduction

43 The investigation of the propagation of natural light in the ocean can be addressed experimentally through in
44 situ measurements and theoretically through numerical simulations of radiative transfer. The understanding of the
45 relationships between the radiometric quantities (i.e., radiance and irradiances) that characterize the light fields
46 within and leaving the ocean and the inherent optical properties (IOPs) of the water column, boundary conditions
47 at the sea surface (i.e., surface illumination conditions and sea state) and at the ocean bottom (i.e., bottom depth
48 and albedo) requires comprehensive datasets of multiple variables acquired over a broad range of environmental
49 conditions. For example, of particular interest are the relationships between the spectral remote-sensing reflectance
50 of the ocean (in sr^{-1}), $R_{rs}(\lambda)$, which is an apparent optical property (AOP) derivable from radiometric quantities,
51 and the seawater IOPs that are directly linked to various seawater constituents because these relationships form
52 the cornerstone of various applications of optical (ocean color) remote sensing. Recent technological developments
53 and broader accessibility of optical in situ instrumentation have led to significant increase in optical datasets
54 collected across diverse oceanic environments and efforts have been undertaken to merge data from various
55 sources within publicly available databases (e.g., Werdell and Bailey, 2005; Valente et al., 2019; Casey et al.,
56 2020). Although the importance of field data collection across diverse environments cannot be overstated, the
57 existing database compilations are subject to certain limitations. In addition to typical measurement errors, it is
58 difficult to ensure consistent data quality and characterization of uncertainties across all merged data because
59 individual datasets are often obtained with different instruments as well as measurement and data processing
60 methods. Also, even the large databases such as [NASA's SeaWiFS Bio-optical Archive and Storage System](https://seabass.gsfc.nasa.gov/)
61 ([SeaBASS, https://seabass.gsfc.nasa.gov/](https://seabass.gsfc.nasa.gov/)) ~~NASA's SeaBASS~~ cannot ensure the balanced representativeness of
62 collected field data in terms of a broad range of optical conditions across diverse ocean environments. In this
63 context, radiative transfer (RT) simulations, which are free of measurement errors, provide a useful tool to generate
64 comprehensive synthetic databases and complement the existing datasets of field measurements in support of
65 studies in ocean optics and optical remote sensing.

66 Over the past decades various radiative transfer models that employ different numerical solution techniques
67 have been developed and used to address a wide range of problems related to optics of natural water bodies (e.g.,
68 Mobley et al., 1993; Mobley, 1994; Stamnes et al., 2017). Since the 1990s the HydroLight code based on invariant
69 imbedding technique (Mobley, 1989; Mobley et al., 1993; Mobley, 1994) has been among the most commonly
70 used radiative transfer models in oceanographic optics. The HydroLight code solves the scalar (i.e., polarization
71 of light is not included) time-independent radiative transfer equation for a horizontally homogeneous water body,
72 in which the inherent optical properties can vary with depth, under given boundary conditions at the surface and
73 bottom of the water body. The inelastic radiative processes within the water column that include Raman scattering
74 by water molecules, fluorescence of chlorophyll-a pigment, and fluorescence of colored dissolved organic matter
75 (CDOM) can be included in HydroLight simulations.

76 The radiative transfer simulations with HydroLight code have proven useful for generating synthetic databases
77 of light field characteristics (i.e., radiance and irradiances) within and leaving the ocean and the AOPs derived
78 from the simulated radiometric quantities for various scenarios of seawater IOPs that provide input to the
79 simulations. In particular, as a result of efforts dedicated to inverse bio-optical algorithms and coordinated under

80 the auspices of the International Ocean Colour Coordinating Group (IOCCG Report, 2006), a widely-used publicly
81 available synthetic database was generated within the spectral range 400 - 800 nm with a 10 nm resolution for
82 clear sky conditions with three different sun zenith angles (0° , 30° , and 60°), a sea surface state corresponding to
83 a wind speed of 5 m s^{-1} , and 500 different IOP combinations driven by chlorophyll-a concentration, Chla, within
84 the surface ocean layer. The input IOP data included the spectral absorption coefficients of phytoplankton, $a_{\text{ph}}(\lambda)$,
85 non-algal particles (also referred to as depigmented or detrital particles ~~which-that~~ can include various types of
86 particles such as organic detritus, mineral particles, heterotrophic bacteria, and depigmented phytoplankton cells),
87 $a_{\text{d}}(\lambda)$, colored dissolved organic matter (CDOM), $a_{\text{g}}(\lambda)$, and the spectral backscattering coefficients of
88 phytoplankton, $b_{\text{b-ph}}(\lambda)$, and non-algal particles, $b_{\text{b-d}}(\lambda)$ (λ represents the wavelength of light in vacuum in units of
89 nm and the IOP coefficients are typically expressed in units of m^{-1}). The output parameters provided by those
90 simulations ~~which-that~~ are available in the public database included the following AOPs: the spectral remote-
91 sensing reflectance, $R_{\text{rs}}(\lambda)$, the remote-sensing reflectance just below the sea surface, $r_{\text{rs}}(\lambda)$, the irradiance
92 reflectance just below the sea surface, $R(z=0^-, \lambda)$, and the diffuse attenuation coefficient for downwelling plane
93 irradiance, $K_{\text{d}}(\lambda, z)$, at the depths $z = 0^-, 5$, and 10 m (where 0^- indicates the depth just beneath the sea surface).

94 Another synthetic database that is publicly available was developed as part of the CoastColour Round Robin
95 project (Nechad et al., 2015). This project was focused on coastal waters and IOPs ~~were driven~~ were described by
96 5000 combinations of Chla, $a_{\text{g}}(\lambda)$, and mass concentration of mineral particles. The HydroLight simulations were
97 run from 350 nm to 900 nm at 5 nm intervals for cloudless sky, three sun zenith angles (0 , 40 , and 60°), and a wind
98 speed of 5 m s^{-1} . The output parameters included in the publicly available database are the water leaving
99 reflectance, $RL_{\text{w}}(\lambda) = \pi R_{\text{rs}}(\lambda)$, $K_{\text{d}}(\lambda)$, the photosynthetically available radiation, PAR , and the euphotic depth, z_{eu} .
100 Most recently, a synthetic database was also developed by the first NASA PACE (Plankton, Aerosols, Cloud, ~~and~~
101 Ocean Ecosystem) Science Team where the ocean contribution to the top of the atmosphere radiances were
102 simulated by HydroLight (Craig et al., 2020). These simulations were performed from 350 to 800 nm with a 5 nm
103 step for a cloudless sky, three sun zenith angles (10° , 30° , and 60°), wind speed of 5 m s^{-1} , and a set of 720 IOP
104 combinations driven by $a_{\text{ph}}(\lambda)$. The publicly available output of these HydroLight simulations is $R_{\text{rs}}(\lambda)$.

105 While these existing synthetic databases have offered valuable information to the ocean color radiometry
106 (OCR) community, especially for the purpose of algorithm development where the ocean AOPs are linked to IOPs,
107 there are several reasons that have motivated the present study aiming at generating a new optical synthetic
108 database. First, the inelastic Raman scattering and fluorescence processes were ignored in the previous RT
109 simulations. These inelastic radiative processes are known to be important for simulating realistic characteristics
110 of light fields within and leaving the ocean, including $R_{\text{rs}}(\lambda)$ that is a primary optical quantity used in ocean color
111 remote sensing. For example, Raman scattering by water molecules may have an important influence on light
112 within and leaving the ocean and AOPs, especially in the green and red parts of the spectrum (e.g., Marshall and
113 Smith, 1990; Stavn, 1993; Sugihara et al., 1984; Westberry et al., 2013). Second, the three synthetic databases
114 described above are based on the use of the spectral pure seawater absorption, $a_{\text{w}}(\lambda)$, and scattering $b_{\text{w}}(\lambda)$,
115 coefficients values as defined by Pope and Fry (1997) and Morel (1974) in the visible part of the spectrum,
116 respectively. However, more recent measurements and theoretical considerations provide new recommendations
117 for spectral values of $a_{\text{w}}(\lambda)$ and $b_{\text{w}}(\lambda)$ (IOCCG Protocol Series, 2018; Zhang and Hu, 2009; Zhang et al., 2009).
118 Third, the ~~suites of IOP data combinations and~~ probability distributions of different IOPs that were used as input
119 to previous RT simulations do not appear to match well with the IOP distributions observed in extensive field

120 datasets or satellite-derived datasets representing the global ocean. This issue may have a biasing effect when the
121 synthetic databases are used to develop the optical algorithms based on the AOP vs. IOP relationships, especially
122 when the underlying goal is to represent a broad range of IOPs encountered within the global ocean, even if the
123 primary interest is in open-ocean pelagic environments. Finally, the previous synthetic databases were developed
124 specifically for OCR-oriented studies and the publicly accessible data generally include only the surface
125 reflectances, $R_{rs}(\lambda)$, $R(\lambda)$, $r_{rs}(\lambda)$, and $K_d(\lambda)$ at selected depths. These databases do not include many of the various
126 output variables obtained from RT simulations, such as the various underwater AOPs, which can be useful in
127 supporting a broader range of studies in ocean optics beyond ocean color remote sensing.

128 In this article, we present a new synthetic optical database generated using RT simulations ~~which-that~~
129 addresses some of the limitations of similar databases developed in the past. First, we describe the development
130 of the synthetic IOP dataset that is required to run RT simulations. The key roles in this development are played
131 by the measured data of phytoplankton absorption coefficient and desired consistency between the probability
132 distributions of synthetic IOPs and the global distributions based on satellite observations. Following this, we
133 describe different configurations of RT simulations ~~which-that~~ were performed with the HydroLight code. The
134 next section is dedicated to consistency between the new optical synthetic database and in situ data, including
135 some previously reported empirical relationships. We provide example illustrations of consistency for both the
136 IOP and AOP data. The closing section summarizes the structure of synthetic database files and provides example
137 illustration of one output radiometric variable, the spectral downwelling plane irradiance, calculated with RT
138 simulations.

139 2 Development of synthetic dataset of seawater inherent optical properties

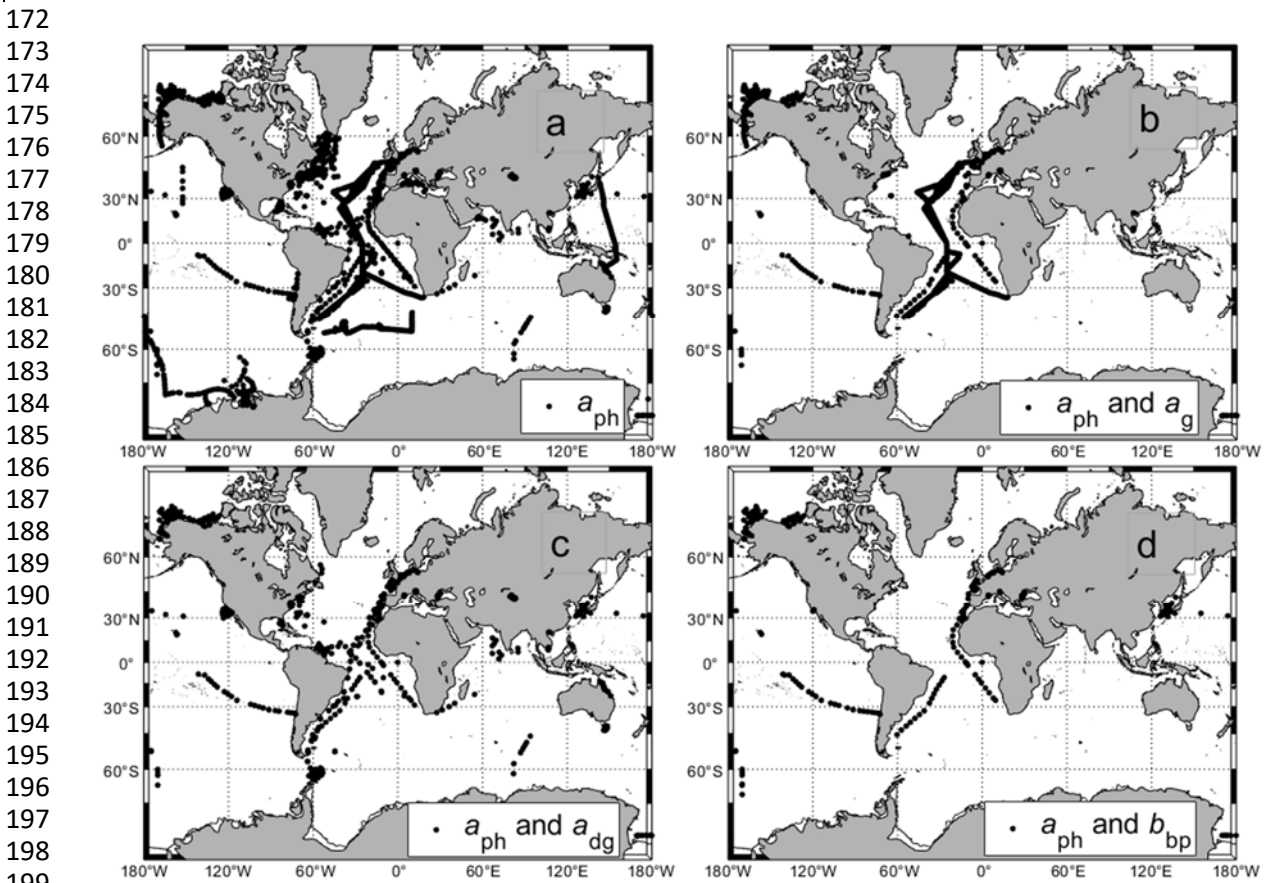
140 [2.1 General overview of methodology](#)

141 The scope of the synthetic database generated with RT simulations and the degree of its representativeness of
142 diverse marine optical environments within the global ocean depend most critically on a dataset of seawater IOPs
143 that are used as input to RT simulations. In the present study, our approach to generate the IOP dataset was driven
144 largely by an underlying goal to obtain the probability distributions of IOPs ~~which-that~~ are generally consistent
145 with the distributions observed in the global ocean dominated by open-ocean pelagic environments. The key IOPs
146 involved in the creation of our IOP dataset include the spectral absorption and backscattering coefficients
147 associated with the main categories of seawater constituents representing suspended particulate matter and CDOM.
148 Specifically, the ~~constituent-specific~~ absorption ~~coefficients of the different constituents~~ ~~IOPs~~ are the spectral
149 absorption coefficients of phytoplankton, $a_{ph}(\lambda)$, non-algal particles, $a_d(\lambda)$, and CDOM, $a_g(\lambda)$. Note that the sum
150 $a_{ph}(\lambda) + a_d(\lambda) = a_p(\lambda)$ represents the particulate absorption coefficient with combined contributions of
151 phytoplankton and non-algal particles, and the sum $a_d(\lambda) + a_g(\lambda) = a_{dg}(\lambda)$ represents the non-phytoplankton
152 absorption coefficient with combined contributions of non-algal particles and CDOM. The ~~constituent-specific~~
153 backscattering ~~coefficient of the different constituents~~ ~~IOPs~~ are the spectral backscattering coefficients of
154 phytoplankton, $b_{b-ph}(\lambda)$, and non-algal particles, $b_{b-d}(\lambda)$, such that the sum $b_{b-ph}(\lambda) + b_{b-d}(\lambda) = b_{bp}(\lambda)$ is the
155 particulate backscattering coefficient.

156 Among these constituent-~~specific~~ IOPs, the phytoplankton absorption coefficient, $a_{ph}(\lambda)$, plays the most
157 fundamental role in the creation of the synthetic dataset of IOPs in this study. The $a_{ph}(\lambda)$ spectra in this dataset
158 were derived from actual measurements of phytoplankton absorption made on near-surface samples collected

159 across diverse oceanic environments. Thus, the $a_{ph}(\lambda)$ data are not “synthetic” in a sense that these data were not
 160 obtained from a modeling approach although some spectral interpolation or extrapolation was applied to measured
 161 data as described in more detail below. In contrast, the remaining four constituent-specific IOPs in the IOP dataset,
 162 i.e., $a_d(\lambda)$, $a_g(\lambda)$, $b_{b-ph}(\lambda)$, and $b_{b-d}(\lambda)$, are “synthetic” in a sense that they are entirely based on calculations using a
 163 modeling approach with some assumptions about the magnitude and spectral behavior of the modeled IOPs.
 164 Importantly, the measured values of $a_{ph}(\lambda)$ were used in the calculations of these IOPs. These calculations are also
 165 described in detail below. Thus, each combination of the five constituent-specific IOPs in the synthetic IOP dataset
 166 consists of the measured $a_{ph}(\lambda)$ and the calculated $a_d(\lambda)$, $a_g(\lambda)$, $b_{b-ph}(\lambda)$, and $b_{b-d}(\lambda)$ where the results of these
 167 calculations depend on the measured $a_{ph}(\lambda)$. As a result of this approach, it would seem justifiable to refer to the
 168 created IOP dataset as a quasi-synthetic dataset. For simplicity, however, we refer to it as the synthetic IOP dataset
 169 while bearing in mind that $a_{ph}(\lambda)$ spectra were derived from measurements.

170
 171 2.2 Description of in situ dataset

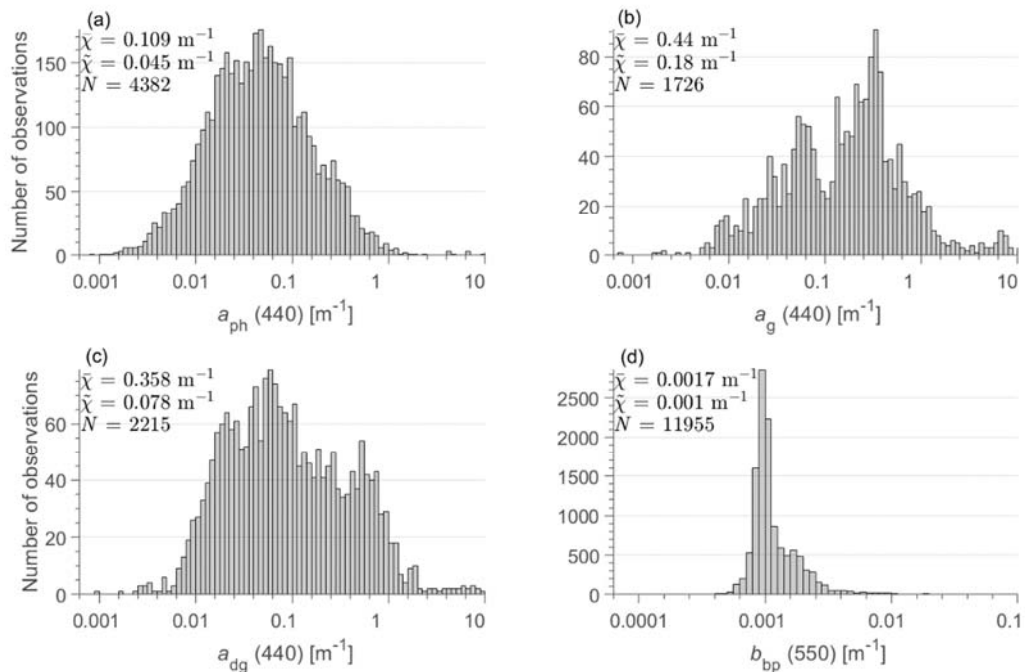


200 Figure 1. Location of oceanographic stations where in situ measurements were collected for (a) $a_{ph}(\lambda)$, the number of
 201 measurements $N = 4382$; (b) $a_{ph}(\lambda)$ and $a_g(\lambda)$, the number of matchup measurements $N = 2206$; (c) $a_{ph}(\lambda)$ and $a_{dg}(\lambda)$, the
 202 number of matchup measurements $N = 813$; and (d) $a_{ph}(\lambda)$ and $b_{bp}(\lambda)$, the number of matchup measurements $N = 775$.

203
 204 Figure 1a depicts the location of oceanographic stations where the near-surface measurements of $a_{ph}(\lambda)$ were
 205 made. As shown, these measurements were collected across diverse open ocean and coastal environments and their
 206 total number is 4382 that constitutes the initial field dataset of $a_{ph}(\lambda)$ considered in this study. Figure 1 also shows
 207 the location of stations where coincident measurements are available for the pairs of IOP coefficients, namely

208 $a_{ph}(\lambda)$ and $a_g(\lambda)$ (Fig. 1b), $a_{ph}(\lambda)$ and $a_{dg}(\lambda)$ (Fig. 1c), and $a_{ph}(\lambda)$ and $b_{bp}(\lambda)$ (Fig. 1d). We recall that while the in
 209 situ data of $a_g(\lambda)$, $a_{dg}(\lambda)$, and $b_{bp}(\lambda)$ were not used in the development of synthetic IOP dataset, they were assembled
 210 for the purpose of comparison with corresponding coefficients that were calculated and included in the synthetic
 211 IOP dataset. Many in situ data of IOP coefficients used in the present study were utilized in previous studies (e.g.,
 212 Babin et al., 2003; Huot et al., 2008; Bricaud et al., 2010; Antoine et al., 2011; Loisel et al., 2018; Aurin et al.,
 213 2018; Reynolds and Stramski, 2019; Stramski et al., 2019; Jorge et al., 2021; Bonelli et al., 2021). Some data are
 214 described in publications devoted to compilation of various datasets (Valente et al., 2019; Casey et al., 2020) and
 215 are included in several databases (e.g., SeaBASS, CoastIOOC, and GOCAD). Many in situ data of IOP coefficients
 216 used in the present study were collected in previous studies (e.g., Reynolds et al., 2001; Babin et al., 2003; Loisel
 217 et al., 2007; Claustre et al., 2008; Huot et al., 2008; Stramski et al., 2008; Lubac et al., 2008; Loisel et al., 2009;
 218 Bricaud et al., 2010; Loisel et al., 2011; Antoine et al., 2011; Neukermans et al., 2012; Uitz et al., 2015;
 219 Neukermans et al., 2016; Reynolds et al., 2016; Aurin et al., 2018; Reynolds and Stramski, 2019; Stramski et al.,
 220 2019). Some data are described in publications devoted to compilation of various datasets (Valente et al., 2019;
 221 Casey et al., 2020) and are included in several databases (e.g., SeaBASS, CoastIOOC, BOUSSOLE, and GOCAD).
 222 As the IOP coefficients in the in situ dataset were measured over a broad range of trophic and environmental
 223 conditions, their spectral values span more than 3 or 4 orders of magnitude. This large dynamic range is illustrated
 224 in terms of probability distributions at selected light wavelengths, i.e., at 440 nm for the constituent absorption
 225 coefficients and 550 nm for the particulate backscattering coefficient (Fig. 2).

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228
229 Figure 2. Histograms and relevant statistical parameters of field measurements of (a) $a_{ph}(440)$, (b) $a_g(440)$, (c) $a_{dg}(440)$, and
 230 (d) $b_{bp}(550)$. N is the number of measurements, and \bar{x} and \tilde{x} are the mean and median values of each IOP, respectively.
 231

232 [2.3 Generation of the dataset of hyperspectral \$a_{ph}\(\lambda\)\$](#)

233 The first task necessary for development of the synthetic IOP dataset was to assemble data of hyperspectral
234 absorption coefficient of phytoplankton, $a_{ph}(\lambda)$, from field measurements collected across diverse open ocean and
235 coastal environments (Fig. 1a, Fig. 2a). These $a_{ph}(\lambda)$ data were obtained with the filter-pad spectrophotometric
236 method as a difference between the measurements of $a_p(\lambda)$ and $a_d(\lambda)$ (Kishino et al., 1985; IOCCG Protocol Series,
237 2018). Historically, most of these measurements were acquired with the transmittance configuration of the filter-
238 pad method and such measurements are included in our dataset. However, some data in our dataset were obtained
239 with the inside integrating-sphere configuration of the filter-pad method, which is superior to the transmittance
240 configuration of measurement (Stramski et al., 2015; IOCCG Protocol Series, 2018).

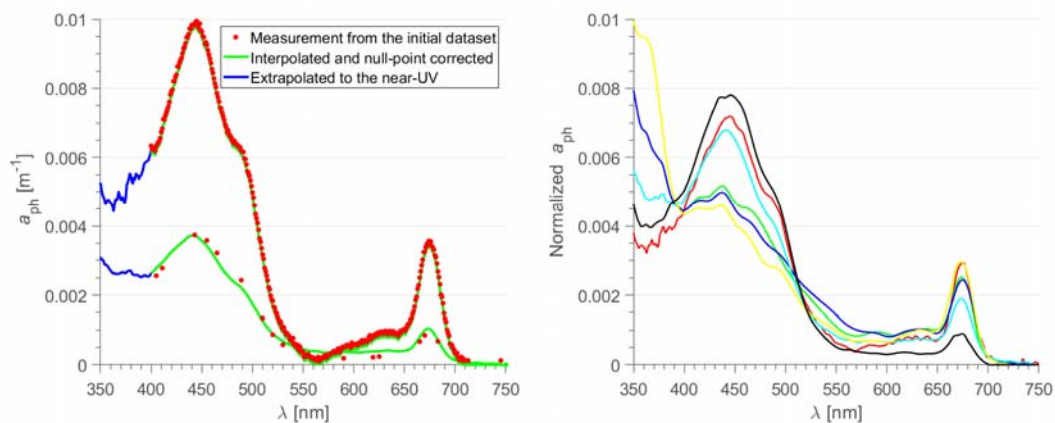
241 A significant portion (23.7%) of the initial dataset of $a_{ph}(\lambda)$ consisting of 4382 measurements covers a spectral
242 range from 400 to 750 nm with high spectral resolution of data reported at 1 nm interval. In some cases, the original
243 measurements extended to near-UV spectral region and/or longer wavelengths in the near-infrared spectral region
244 (800 or 850 nm). The data beyond 750 nm are not used in this study because our RT simulations target the spectral
245 range from 350 to 750 nm. It is notable that the absorption measurements of marine particles and phytoplankton
246 are generally unavailable or are not reported in the UV because of increased methodological challenges and
247 uncertainties in this spectral region (Stramski et al., 2015; IOCCG Protocol Series, 2018; Kostakis et al., 2021).
248 As a result, only a relatively small fraction of $a_{ph}(\lambda)$ measurements in the initial dataset were reported in the near-
249 UV region. In addition, the initial dataset included a relatively large fraction of $a_{ph}(\lambda)$ measurements that were
250 reported at wavelength intervals larger than 1 nm. These lower resolution data (hereafter referred to as
251 multispectral) ranged from a small wavelength interval of 2 nm to data reported at more limited number of
252 wavelengths (as small as <10) within the visible spectral range. It is likely that the multispectral data available
253 from some data sources that we used in this study were originally measured at higher spectral resolution but
254 eventually were reported only for some selected wavelengths, such as those corresponding to spectral bands
255 available on satellite ocean color sensors.

256 The first objective of the analysis of $a_{ph}(\lambda)$ was to consider the initial $a_{ph}(\lambda)$ dataset within the 400–750 nm
257 range and convert the measurements that were reported at lower spectral resolution to uniformly hyperspectral
258 data at 1 nm interval. In this analysis, all measurements originally available at 1 nm interval were considered to
259 provide reference spectral shape functions of $a_{ph}(\lambda)$. The originally multispectral data of $a_{ph}(\lambda)$ were converted to
260 hyperspectral data using ~~a couple of several~~ different approaches depending on the spectral features of lower
261 resolution data. One approach utilized the reference spectral shape functions of $a_{ph}(\lambda)$ and was applied to
262 multispectral $a_{ph}(\lambda)$ data if they were reported at fewer wavelengths than 100. In this case, a given multispectral
263 spectrum of $a_{ph}(\lambda)$ was converted to hyperspectral spectrum using a specific hyperspectral measurement that
264 exhibited the highest correlation with the multispectral measurement under consideration. The correlation
265 coefficient was calculated using the spectral data available at common wavelengths of considered pair of spectra.
266 A necessary condition to proceed with a conversion of a given multispectral spectrum to hyperspectral spectrum
267 was a correlation coefficient of 0.95 or higher. If this condition was satisfied, the multispectral data were converted
268 to hyperspectral data so that the created hyperspectral spectrum maintained the magnitude of multispectral
269 measurement in the range of 440–450 nm and had the spectral shape of the reference hyperspectral measurement.
270 An alternative approach to convert multispectral data to hyperspectral data involved a linear interpolation of
271 multispectral data. This approach was used when the multispectral data were reported at relatively small
272 wavelength intervals (at least 100 spectral data available between 400 and 750 nm) or when the correlational

273 analysis described above did not yield the correlation coefficient of 0.95 or higher (5.2% of the multispectral data).
 274 The originally multispectral spectra that-which did not include data below 450 nm or fell into the category of data
 275 subject to linear interpolation but had no data above 700 nm were rejected from further analysis. For all
 276 hyperspectral spectra that passed the above-described analysis and criteria (i.e., 2204 spectra that included both
 277 the 593 original hyperspectral measurements and 1611 hyperspectral spectra created from multispectral data), the
 278 null-point correction was applied by subtracting the average value of $a_{ph}(\lambda)$ in the 745–750 nm range from all
 279 spectral values in the 400–750 nm range.

280 The next step of analysis was to extend all null-point corrected spectra of $a_{ph}(\lambda)$ that cover the 400–750 nm
 281 range into the UV spectral region. The primary focus was on the 350–400 nm range because our RT simulations
 282 were designed to provide output results in the 350–750 nm range. For this purpose we used a separate subset of
 283 reference hyperspectral measurements of $a_{ph}(\lambda)$ that includes the near-UV spectral region. This reference subset
 284 of data consisted of 233 measurements collected across bio-optically diverse marine environments in the Pacific
 285 and Atlantic Oceans and western Arctic seas. The majority of these 233 spectra (170) were collected with the
 286 inside integrating-sphere configuration of filter-pad method, while the remaining 63 measurements were done
 287 using either the transmittance or transmittance-reflectance filter-pad configuration (Zheng et al., 2014). A
 288 correlational analysis was applied to pairs of spectra, each consisting of a spectrum covering the 400–750 nm range
 289 and a reference spectrum covering the 350–750 nm range. The correlation coefficient was calculated using data at
 290 common wavelengths from the 400–750 nm range. The reference spectrum that yielded the highest correlation
 291 with the investigated 400–750 nm spectrum was selected as a basis for extrapolation of the investigated spectrum
 292 into the 350–400 nm range. This extrapolation ensured that a given investigated spectrum maintained its magnitude
 293 at 400 nm and its extrapolated near-UV portion had the spectral shape of the selected reference spectrum. The final
 294 aspect of extrapolation in the UV is related to the spectral range 300–350 nm. The rationale for IOP data extending
 295 to 300 nm is to ensure that the results of RT simulations that start at 350 nm account for possible effects of Raman
 296 scattering by water molecules in the UV spectral region. Therefore, for the 300–350 nm range we simply assumed
 297 that $a_{ph}(\lambda)$ in this range is equal to $a_{ph}(350)$. The limitation associated with this assumption is not considered to be
 298 serious given the limited role of the 300–350 nm range in the RT simulations and weak Raman scattering effects
 299 in UV spectral region. Example spectra of $a_{ph}(\lambda)$ in the 350–750 nm range from contrasting marine environments
 300 are presented in Fig. 3. These examples show significant variation in both the magnitude and spectral shape of
 301 $a_{ph}(\lambda)$.

302



303
 304

305 Figure 3. (a) Two example spectra of $a_{ph}(\lambda)$ from contrasting oceanic environments. For each example $a_{ph}(\lambda)$, two spectra are
306 displayed, namely the measurement from the initial $a_{ph}(\lambda)$ dataset shown at the original wavelength intervals (red points) and
307 the spectrum after interpolation to 1 nm intervals (if required) and null-point correction (continuous lines). The UV portion of
308 the latter was obtained by extrapolation based on reference data in the UV (see text for details). (b) Example of normalized
309 $a_{ph}(\lambda)$ spectra illustrating the variability of the spectral shape of the $a_{ph}(\lambda)$ database. These spectra have been normalized to
310 their integral.

311

312 [2.4 Generation of the complete IOP dataset](#)

313 In the next step of analysis, the subset of 2204 $a_{ph}(\lambda)$ spectra that was created from the initial $a_{ph}(\lambda)$ dataset as
314 described above was subject to additional modifications to ensure that the final $a_{ph}(\lambda)$ dataset is characterized by
315 the probability distribution that resembles the distribution representative of the global ocean. This process and
316 background information on the motivation for such adjustments in the probability distribution are described below.

317 When the end goal is to achieve a high degree of representativeness of global ocean like in this study, the
318 process of assembling in situ datasets of IOPs is unavoidably subject to limitations, even if relatively large amount
319 of data from many field experiments and cruises are considered. This is mainly because the global ocean is
320 dominated by vast areas of open-ocean pelagic environments and the amount of IOP data collected in these
321 environments is disproportionally limited compared to amount of data collected in coastal regions that represent a
322 relatively small portion of the global ocean. Thus, the probability distributions based on in situ datasets, such as
323 those presented in Fig. 2, are expected to deviate from the probability distributions representative of the global
324 ocean. In particular, the maxima of probability distributions and the measures of central tendency, such as the
325 median and mean values, obtained from compilations of relatively large amount of in situ IOP data (such as in Fig.
326 2) are expected to be shifted to larger values compared to actual global distributions because the IOPs exhibit a
327 general tendency of higher values in coastal regions compared to open ocean environments. While this issue has
328 been recognized, it has not been addressed or resolved in various studies that focus on global ocean color
329 applications. For example, the current global ocean color algorithms for estimating chlorophyll-a concentration
330 (Chla) are based on relatively large amount of in situ data whose probability distribution is shifted significantly to
331 higher Chla compared with the global Chla distribution (O'Reilly and Werdell, 2019). Similarly, in the
332 development of previous synthetic optical databases with RT simulations (e.g., IOCCG Report, 2006), no special
333 attempt was made to ensure consistency between the probability distributions of input IOP data and the
334 distributions expected for global ocean. In the recent development of refined global ocean color algorithms for
335 estimating the concentration of particulate organic carbon (POC), the in situ dataset was assembled with a goal to
336 achieve reasonable consistency with a global POC distribution (Stramski et al., 2022). This goal was, however,
337 achieved at the expense of significant reduction in the amount of accepted in situ data compared to the size of
338 overall pool of available in situ data.

339 In this study our goal was to create a relatively large synthetic IOP dataset based on the initial dataset of
340 several thousand measurements of spectral $a_{ph}(\lambda)$, so that the probability distributions of IOPs in the final synthetic
341 dataset are reasonably consistent with the expected distributions representative of the global ocean. As described
342 above, the initial field dataset in support of this process consisted of 4382 spectra of $a_{ph}(\lambda)$ and this number was
343 further reduced to 2204 spectra that were accepted as a result of analysis and some criteria applied to the initial
344 dataset. This reduced dataset of accepted $a_{ph}(\lambda)$ spectra was then further modified to ensure that the final
345 probability distribution of $a_{ph}(440)$ resembles the global distribution of $a_{ph}(440)$. The global probability

346 distribution of $a_{\text{ph}}(440)$ was estimated using retrievals of $a_{\text{ph}}(440)$ from satellite ocean color data. Specifically, we
 347 used global satellite observations made with the ocean color sensor OLCI (Ocean and Land Colour Instrument)
 348 deployed on the Sentinel-3 mission (Donlon et al., 2012) from the period December 1, 2020 through November
 349 30, 2021. The weekly data product of remote-sensing reflectance $R_{\text{rs}}(\lambda)$ at 4 km² spatial resolution was used as
 350 input to the 3-step semi-analytical algorithm (3SAA) to derive $a_{\text{ph}}(443)$ as described in Jorge et al (2021). The
 351 $a_{\text{dg}}(443)$ and $b_{\text{bp}}(\lambda)$ coefficients were also derived from this algorithm. In general, the 3SAA first derives the diffuse
 352 attenuation coefficient for downwelling plane irradiance averaged within the surface layer down to the first
 353 attenuation depth, $\langle K_{\text{d}}(\lambda) \rangle_1$, from $R_{\text{rs}}(\lambda)$, and then utilizes the inverse model LS2 (Loisel et al., 2018) to derive the
 354 total absorption, $a(\lambda)$, and backscattering, $b_{\text{b}}(\lambda)$, coefficients from $R_{\text{rs}}(\lambda)$ and $\langle K_{\text{d}}(\lambda) \rangle_1$. After subtracting the pure
 355 seawater contributions, the non-water absorption, $a_{\text{nw}}(\lambda)$, and the particulate backscattering, $b_{\text{bp}}(\lambda)$, coefficients
 356 are obtained. Finally, $a_{\text{ph}}(\lambda)$ and $a_{\text{dg}}(\lambda)$ are derived from $a_{\text{nw}}(\lambda)$ using an optimization algorithm of Zhang et al.
 357 (2015) with modifications that account for differences in optical water types defined in terms of different spectral
 358 shapes of $R_{\text{rs}}(\lambda)$ (Mélin and Vantrepotte, 2015). While the original classification of Mélin and Vantrepotte (2015)
 359 includes 16 optical water classes (OWC), the derivation of $a_{\text{ph}}(\lambda)$ and $a_{\text{dg}}(\lambda)$ from the 3SAA additionally included
 360 a 17th OWC to improve the representation of ultraoligotrophic waters such as those found in the South Pacific
 361 Gyre (Morel et al., 2007; Claustre et al., 2008; Stramski et al., 2008) and in some areas of the Mediterranean Sea
 362 in summer (Loisel et al., 2011). This 17th OWC is described in Jorge et al. (2021).

363 The 3SAA does not yield the separate contributions of CDOM, $a_{\text{g}}(\lambda)$, and non-algal particles, $a_{\text{d}}(\lambda)$, to the
 364 overall non-phytoplankton absorption coefficient, $a_{\text{dg}}(\lambda)$. Therefore, we also used another semi-analytical model
 365 (CDOM-KD2) described in Bonelli et al. (2021) to estimate $a_{\text{g}}(443)$ from OLCI-derived $R_{\text{rs}}(\lambda)$. Having $a_{\text{dg}}(\lambda)$ from
 366 the 3SAA and $a_{\text{g}}(\lambda)$ from the CDOM-KD2, the non-algal particulate absorption, $a_{\text{d}}(\lambda)$, was obtained as a difference
 367 $a_{\text{dg}}(\lambda) - a_{\text{g}}(\lambda)$. As a result of this analysis, we obtained a dataset of satellite-derived constituent absorption
 368 coefficients, $a_{\text{ph}}(443)$, $a_{\text{g}}(443)$, $a_{\text{d}}(443)$, and $a_{\text{dg}}(443)$, as well as the particulate backscattering coefficient, $b_{\text{bp}}(550)$,
 369 where we focused on the spectral band near 440 nm for absorption and 550 nm for backscattering.

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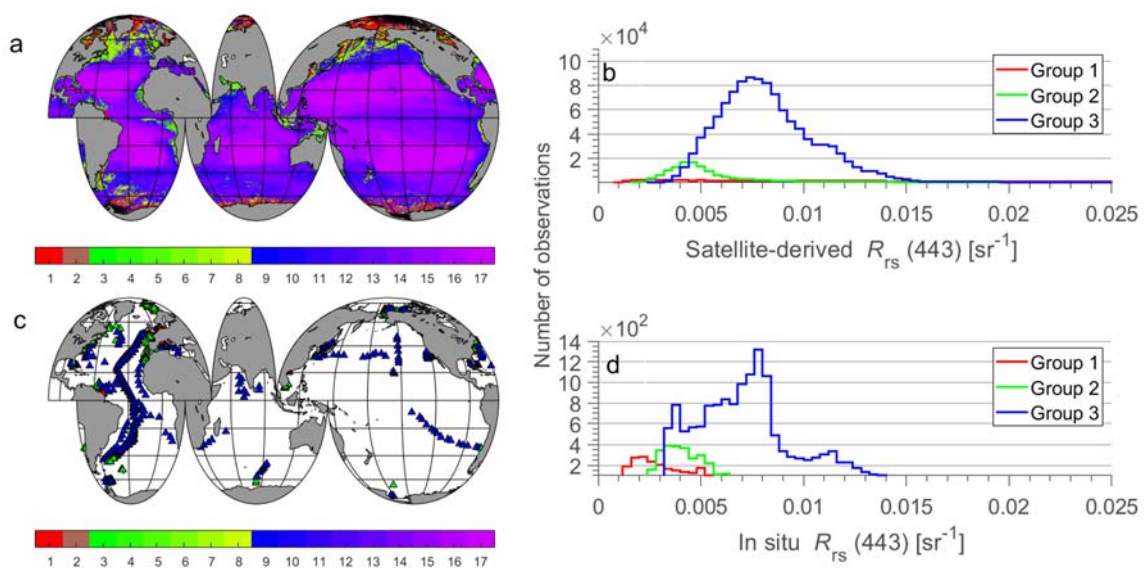
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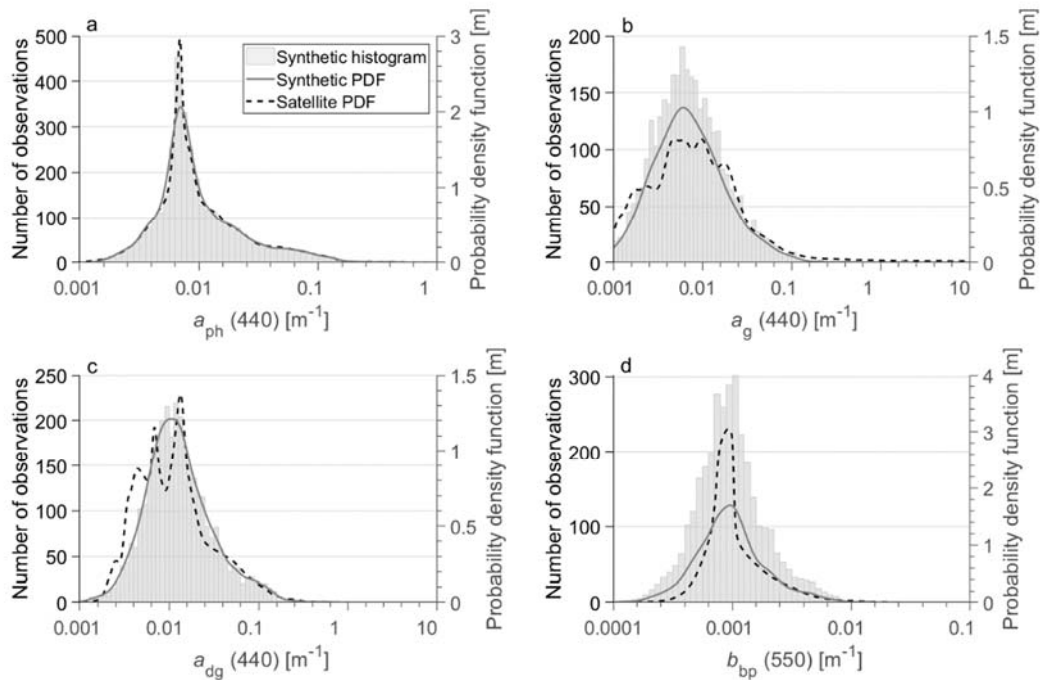
384 Figure 4. (a) Global map illustrating the distribution of seventeen optical water classes estimated from monthly $R_{\text{rs}}(\lambda)$ values
 385 derived from satellite observations with ocean color sensor OLCI from December 2020 through November 2021 (weekly

386 products at 4km²). The color bar scale refers to optical water classes. (b) Histogram of OLCI-derived $R_{rs}(443)$ for the three
 387 optical water groups (see text for details). (c) Location of oceanographic stations where in situ measurements of $R_{rs}(\lambda)$ were
 388 collected and used to analyze the consistency of the synthetic dataset with field measurements. (d) Histograms of in situ
 389 measurements of $R_{rs}(443)$ for the three optical water groups.

390

391 For illustrative purposes Fig. 4a depicts the spatial distribution of 17 optical water classes (OWCs) over the
 392 global ocean obtained from satellite OLCI data following the methodology of Mélin and Vantrepotte (2015). For
 393 further illustrative purposes, these 17 OWCs were grouped into 3 optical water groups (OWGs). Group 1 consists
 394 of OWC1 and OWC2 which are characterized by high water turbidity such as in coastal areas affected by discharge
 395 from large rivers. Although the focus of this study is to create the synthetic datasets representative primarily of
 396 open ocean and moderately turbid coastal waters, an explicit identification of Group 1 data that represent very
 397 turbid waters is of interest for comparisons with the database developed specifically for coastal waters by Nechad
 398 et al. (2015). The second OWG, Group 2, includes 6 OWCs from OWC3 through OWC8. This group represents
 399 mainly productive waters in both coastal and open ocean environments, such as those encountered in the North
 400 Atlantic during a period of phytoplankton bloom (Levy et al., 2005). Finally, Group 3 included the remaining 9
 401 OWCs from OWC9 through OWC17. These water types are observed mainly in mesotrophic and oligotrophic
 402 regions of the global ocean. Based on this classification, 79.6% of OLCI water pixels in Fig. 4a belong to Group
 403 3, 10.8% to Group 2, and 9.6% to Group 1. The histograms of OLCI-derived $R_{rs}(443)$ associated with these three
 404 groups of data are shown in Fig. 4b. For comparative purposes we also assembled a dataset of in situ measurements
 405 of $R_{rs}(\lambda)$, which were collected at various locations within the global ocean (Fig. 4c). The histograms of in situ
 406 $R_{rs}(443)$ associated with Groups 1, 2, and 3 are depicted in Fig. 4d, which show a similar pattern to that in Fig. 4b.
 407 For in situ dataset of $R_{rs}(\lambda)$, 69.2% of data belong to Group 3, 15.7% to Group 2, and 15.1% to Group 1.

408



409

410 Figure 5. Histograms showing the distribution of the synthetic IOP data used in the present study. The synthetic and satellite-
 411 derived probability density functions (PDFs) for each IOP are represented by the solid and dashed curves, respectively.

412

413 The probability density function (PDF) of global satellite-derived $a_{ph}(440)$, $a_g(440)$, $a_{dg}(440)$, and $b_{bp}(550)$ are
414 depicted in Fig. 5. We note that we refer here to satellite-derived absorption coefficients at 440 nm although they
415 were derived from OLCI reflectances at 443 nm, which is a minor difference that is inconsequential for the purpose
416 of this study. The comparison of Fig. 2a and Fig. 5a indicates that the distribution of measured $a_{ph}(440)$ from our
417 initial field dataset (Fig. 2a) is shifted towards higher values compared to the global distribution of satellite-derived
418 $a_{ph}(440)$ (Fig. 5a). The probability distribution of reduced dataset of measured $a_{ph}(440)$ ($N = 2204$) that was created
419 from the initial field dataset of $a_{ph}(\lambda)$ show similar deviations from the global distribution (not shown). Thus, to
420 create the final dataset of $a_{ph}(\lambda)$ that has the probability distribution of $a_{ph}(440)$ consistent with the global satellite-
421 derived distribution, we adjusted the number of $a_{ph}(440)$ measurements in each bin of the histogram of the reduced
422 dataset either by removing the measurements from any given bin or adding the measurements to this bin. The
423 removal or addition of $a_{ph}(440)$ measurements associated with any given bin was done by subjecting all $a_{ph}(440)$
424 measurements originally contained within a given bin to random selection. Specifically, in the case of addition the
425 randomly selected $a_{ph}(440)$ was added as a replicate of $a_{ph}(440)$ to a given bin. In the case of removal, the randomly
426 selected $a_{ph}(440)$ was simply removed from a given bin. As a result of this process we obtained a modified
427 distribution of measured $a_{ph}(440)$ ~~which that~~ is fairly consistent with the satellite-derived distribution of $a_{ph}(440)$.
428 Both the modified histogram and the corresponding modified PDF of measured $a_{ph}(440)$ are depicted in Fig. 5a
429 for comparison with the global satellite-derived distribution. In total, this modified distribution consists of 3320
430 measurements of $a_{ph}(440)$ and, obviously, each of these measurements at 440 nm has an associated full spectrum
431 of $a_{ph}(\lambda)$ values between 300 and 750 nm. These 3320 spectra of $a_{ph}(\lambda)$ represent one IOP component of the final
432 synthetic IOP dataset.

433 The full synthetic IOP dataset created in this study consists of 3320 combinations of measured $a_{ph}(\lambda)$ and
434 synthetically-generated $a_d(\lambda)$, $a_g(\lambda)$, $b_{b-ph}(\lambda)$, and $b_{b-d}(\lambda)$. Below is a description of calculations of $a_g(\lambda)$, $a_d(\lambda)$, b_{b-}
435 $ph(\lambda)$, and $b_{b-d}(\lambda)$. We note that all IOP coefficients are expressed in units of $[m^{-1}]$ and the light wavelength is in
436 units of $[nm]$.

437 The four IOP coefficients were calculated using a similar methodology to that applied in previous studies
438 aiming at generation of synthetic ocean optical databases (IOCCG Report, 2006; Craig et al., 2020). Specifically,
439 we used the measured values of $a_{ph}(440)$ as the main driver of calculations of $a_g(\lambda)$, $a_d(\lambda)$, $b_{b-ph}(\lambda)$, and $b_{b-d}(\lambda)$.
440 Thus, the variability in the measured $a_{ph}(440)$, as depicted by the probability distribution of measured $a_{ph}(440)$ in
441 Fig. 5a, is the main source of variability in these four co-existing IOP coefficients. It is notable that the replicate
442 values of $a_{ph}(440)$ present within any given bin of the $a_{ph}(440)$ distribution result in the generation of different
443 values of the four IOP coefficients because the formulas involved in these calculations contain random parameters.
444 The coupling between $a_{ph}(440)$ and CDOM absorption coefficient was defined as:

$$445 \quad a_g(440) = 10^{(P_1 + \gamma)} \quad (1)$$

446 where P_1 is a parameter related to $a_{ph}(440)$ and γ is randomly selected from a predetermined range of values (Table
447 1). The spectral values of $a_g(\lambda)$ are subsequently determined from:

$$448 \quad a_g(\lambda) = a_g(440) e^{-S_g(\lambda-440)} \quad (2)$$

449 where the spectral slope parameter, S_g in units of $[nm^{-1}]$, is randomly selected from a predetermined range of values
450 (Table 1). The absorption coefficient of non-algal particles was modeled in a similar fashion:

$$451 \quad a_d(440) = P_2 a_{ph}(440) \quad (3)$$

$$a_d(\lambda) = a_d(440) e^{-S_d(\lambda-440)} \quad (4)$$

where P_2 is a parameter related to $a_{ph}(440)$ and the spectral slope parameter S_d [nm^{-1}] is randomly selected from a predetermined range of values (Table 1). The parameterizations of P_1 and P_2 were chosen to match relationships observed with the in situ dataset assembled in this study.

The particulate backscattering is not modeled in terms of the single coefficient, $b_{bp}(\lambda)$, but instead as separate contributions by phytoplankton, $b_{b-ph}(\lambda)$, and non-algal particles, $b_{b-d}(\lambda)$, so that their sum yields $b_{bp}(\lambda)$. In order to calculate $b_{b-ph}(\lambda)$, first the formula that couples $a_{ph}(440)$ with the beam attenuation coefficient of phytoplankton at 550 nm, $c_{ph}(550)$, is used:

$$c_{ph}(550) = P_3 \text{Chla}^{0.57} = P_3 \left[\frac{a_{ph}(440)}{0.05582} \right]^{0.57} \quad (5)$$

where Chla is the concentration of chlorophyll-a in units of [mg m^{-3}], 0.05582 [$\text{m}^2 \text{ (mg}^{-1} \text{ Chla)}$] is the value of chlorophyll-specific absorption coefficient of phytoplankton at 440 nm, $a_{ph}^*(440)$; (Maritorena et al., 2002), and P_3 is a parameter with a randomly selected value from a predetermined range (Table 1). The exponent value of 0.57 is based on the study of Voss (1992). Subsequently, the spectral values of phytoplankton beam attenuation coefficient are calculated from:

$$c_{ph}(\lambda) = c_{ph}(550) \left(\frac{550}{\lambda} \right)^{S_{c-ph}} \quad (6)$$

where the spectral slope parameter, S_{c-ph} [dimensionless], is calculated using both $a_{ph}(440)$ and a random number generator (Table 1). Next, the spectral scattering coefficient of phytoplankton is determined:

$$b_{ph}(\lambda) = c_{ph}(\lambda) - a_{ph}(\lambda) \quad (7)$$

where the spectral values of $a_{ph}(\lambda)$ are from the same measured spectrum as the value of $a_{ph}(440)$ in Eq. (5). Finally, the spectral backscattering coefficient of phytoplankton is calculated from:

$$b_{b-ph}(\lambda) = 0.01 b_{ph}(\lambda) \quad (8)$$

where 0.01 is the value of backscattering ratio of phytoplankton, \tilde{b}_{b-ph} , assumed to be constant and independent of light wavelength (IOCCG, 2006; Loisel et al., 2007; Whitmire et al., 2010). We note that $b_{b-ph}(\lambda)$ is not required as input to our radiative transfer simulations but $b_{ph}(\lambda)$ is needed.

To calculate the backscattering coefficient of non-algal particles, $b_{b-d}(\lambda)$, the phytoplankton absorption at 440 nm is first coupled with the scattering coefficient of non-algal particles at 550 nm, $b_d(550)$, using the following relationship:

$$b_d(550) = P_4 \text{Chla}^{0.766} = P_4 \left[\frac{a_{ph}(440)}{0.05582} \right]^{0.766} \quad (9)$$

where the parameter P_4 is randomly selected from a predetermined range (Table 1) and the value of 0.05582 is $a_{ph}^*(440)$ as explained in relation to Eq. (5). The exponent value of 0.766 is based on the study of Loisel et al. (1998). Then, the spectral values of non-algal scattering coefficient are calculated from:

$$b_d(\lambda) = b_d(550) \left(\frac{550}{\lambda} \right)^{S_{b-d}} \quad (10)$$

where the spectral slope parameter, S_{b-d} [dimensionless], is calculated using both $a_{ph}(440)$ and a random number generator (Table 1). In the final step, the spectral backscattering coefficient of non-algal particles is calculated as:

$$b_{b-d}(\lambda) = 0.018 b_d(\lambda) \quad (11)$$

where the constant 0.018 is the backscattering ratio of non-algal particles, \tilde{b}_{b-d} . This value was proposed by Mobley (1994) and was derived by averaging three particle phase functions measured in oceanic waters by

489 Petzold (1972). Again, we note that $b_{b-d}(\lambda)$ is not required as input to radiative transfer simulations but $b_d(\lambda)$ is
490 needed. The spectral slope of $b_{bp}(\lambda)$, γ , where $b_{bp}(\lambda)$ is obtained as the sum of $b_{b-ph}(\lambda)$ and $b_{b-d}(\lambda)$, has a mean and
491 standard deviation of 1.10 ± 0.34 , and exhibits a decreasing trend-from oligotrophic -(where γ is around -2) to
492 eutrophic waters (where the $b_{bp}(\lambda)$ spectrum is nearly flat). These results are in good agreement with previous
493 studies (Morel and Maritorena, 2001; Loisel et al., 2006; Antoine et al., 2011).

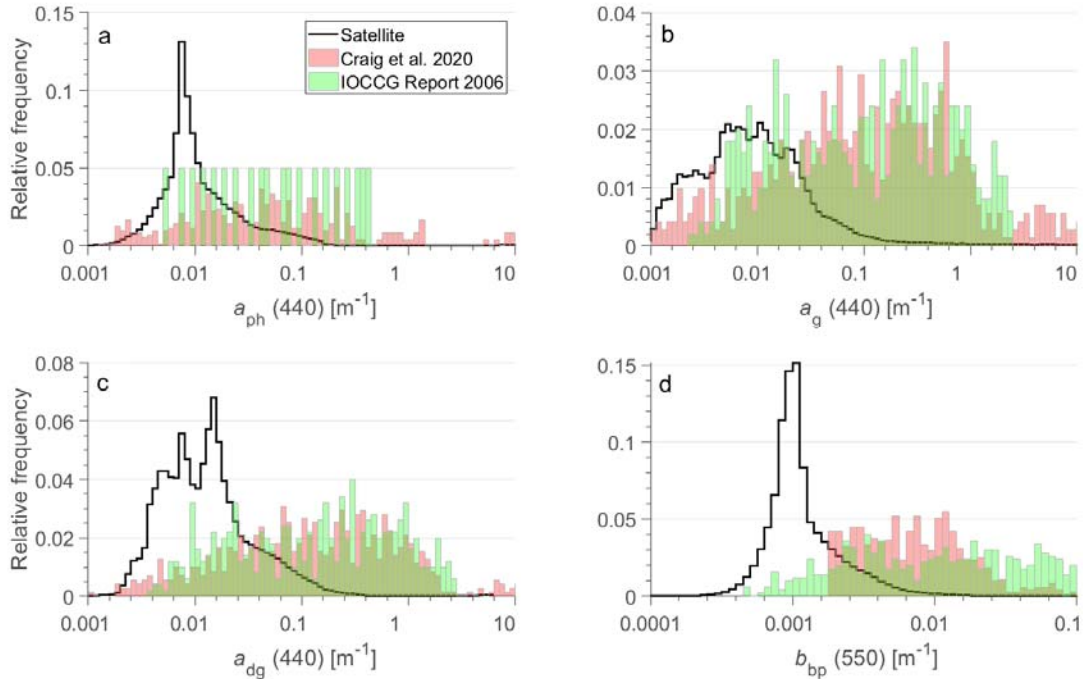
497 Table 1: Symbols of variables, mathematical expressions, and corresponding equations in the text of the paper.
498 $rng(0,1)$ is a random number between 0 and 1.

Symbols	Mathematical expression	Equation	Reference
P_1	$0.79 \log_{10} [a_{ph}(440)] - 0.37$	(1)	This study
γ	$-0.2 + 0.3 \text{rng}(0,1)$	(1)	This study
S_g	$(0.02 - 0.01) \text{rng}(0,1) + 0.01$	(2)	IOCCG (2006)
P_2	$(0.1 + \text{rng}(0,0.9))a_{ph}(440)$	(3)	This study
S_d	$(0.015 - 0.007) \text{rng}(0,1) + 0.007$	(4)	IOCCG (2006)
P_3	$(0.3 - 0.03) \text{rng}(0,1) + 0.03$	(5)	Based on IOCCG (2006)
S_{c-ph}	$-0.4 + \frac{1.6 + 1.2 \text{rng}(0,1)}{1 + [\frac{a_{ph}(440)}{0.05582}]^{0.5}}$	(6)	IOCCG (2006)
P_4	$(0.16668 - 0.016668) \text{rng}(0,1) + 0.016668$	(9)	Based on IOCCG (2006)
S_{b-d}	$-0.5 + \frac{2 + 1.2 \text{rng}(0,1)}{1 + [\frac{a_{ph}(440)}{0.05582}]^{0.5}}$	(10)	IOCCG (2006)

499
500
501 The variability of measured $a_{ph}(440)$ illustrated in Fig. 5a along with the dynamic range of parameters P_1 , P_2 ,
502 P_3 , P_4 , the spectral slopes S_g , S_d , S_{c-ph} , and S_{b-d} , and the degree of randomness in the selection of these parameters
503 for any given value of $a_{ph}(440)$ that initiates the process of calculating $a_g(\lambda)$, $a_d(\lambda)$, $b_{b-ph}(\lambda)$, and $b_{b-d}(\lambda)$, resulted in
504 the generation of synthetic dataset of these IOP coefficients that cover a wide dynamic range consistent with in
505 situ and satellite observations over the global ocean. Figure 5b,c,d compares the probability distributions of
506 satellite-derived $a_g(440)$, $a_{dg}(440)$, and $b_{bp}(550)$ with the distribution of these coefficients from the final synthetic
507 IOP dataset. This comparison supports the general consistency of the distributions of these IOP coefficients, which
508 is in line with the desired consistency achieved for $a_{ph}(440)$ (Fig. 5a) as discussed earlier in this section. It is also
509 noteworthy that in contrast to this newly created synthetic IOP dataset, the previous synthetic datasets exhibit
510 significant differences between the probability distributions of synthetic IOPs and global distributions based on
511 satellite observations (Fig. 6).

512 Overall, the above-described synthetic IOP dataset includes 3320 scenarios of non-water IOPs, i.e., IOPs
513 associated with variable contributions of phytoplankton, non-algal particles, and CDOM to optical properties of
514 seawater. In addition to the non-water absorption coefficients, $a_{ph}(\lambda)$, $a_d(\lambda)$, and $a_g(\lambda)$, as well as the non-water
515 scattering coefficients, $b_{ph}(\lambda)$ and $b_d(\lambda)$, the radiative transfer simulations required input of scattering phase
516 functions of particles, specifically for phytoplankton and non-algal particles. We assumed the particulate phase

517 functions proposed by Fournier-Forand (1994) with the backscattering ratio $\tilde{b}_{b-ph} = 0.01$ for phytoplankton and
 518 $\tilde{b}_{b-d} = 0.018$ for non-algal particles. Note that while the backscattering ratios are assumed to be spectrally constant,
 519 the phase functions vary with light wavelength because of spectral variations of $b_{ph}(\lambda)$ and $b_d(\lambda)$. All IOP data in
 520 the final synthetic IOP dataset cover the spectral range from 300 to 750 nm with a 5 nm interval. This wavelength
 521 interval is consistent with the intended output of our radiative transfer simulations.
 522



523
 524 Figure 6. Histograms showing the distribution of IOPs from the synthetic datasets of the IOCCG Report (2006) and Craig et
 525 al. (2020) in the green and pink, respectively. The IOP distributions estimated from satellite ocean color observations with
 526 OLCI sensor over the global ocean are represented by the black line.

527
 528 The radiative transfer simulations also required input of the absorption and scattering properties of pure
 529 seawater. For the spectral absorption coefficient of pure seawater, $a_w(\lambda)$, we used the values recommended in
 530 IOCCG Protocol Series (2018). This recommendation includes the values from Jonasz and Fournier (2007) in the
 531 spectral range 300-330 nm, Morel et al. (2007) in the 340-415 nm range, Pope and Fry (1997) in the 420-725 nm
 532 range, and Kou et al. (1993) in the 730-750 nm range. The spectral volume scattering function of pure seawater
 533 (from which the spectral scattering coefficient and scattering phase function can be obtained) was calculated
 534 following Zhang et al. (2009) assuming water temperature of 18°C and salinity of 35‰. The temperature of 18°C
 535 is consistent with the mean sea surface temperature (SST) calculated from the monthly global NOAAv2 SST
 536 database at 1° spatial resolution from December 1991 through November 2021 (Jérôme Vialard, personal
 537 communication, <https://www.psl.noaa.gov/data/gridded/data.noaa.oisst.v2.html>). The salinity of 35‰ is also
 538 consistent with the global surface average (Durack et al., 2013).

539 3 Radiative transfer simulations

540 The IOP dataset described in section 2, which includes 3320 combinations of non-water IOPs, provided the
 541 key input to radiative transfer (RT) simulations ~~which-that~~ were performed with the HydroLight v5.0 radiative

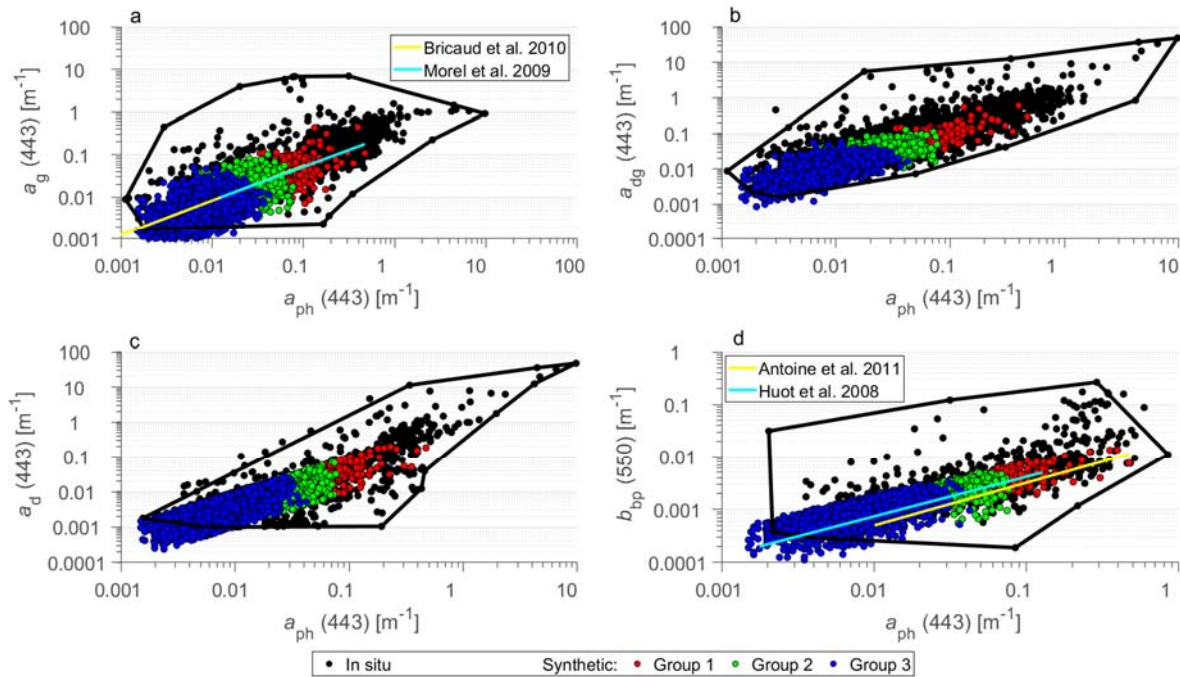
542 transfer code (Mobley and Sundman, 2008). All RT simulations were run assuming vertically homogeneous IOPs
543 within the water column and infinitely deep ocean, i.e., no effect of seafloor on light field within the water column.
544 For all simulations the computed radiometric and AOP variables were saved into the output data files at 10 cm
545 depth intervals between the ocean surface and the 1 m depth, and at 1 m intervals between the 1 m and 50 m depth.
546 Thus, the primary focus of our RT simulations is on the ocean surface layer that can potentially contribute to light
547 leaving the ocean with significance to remote sensing with spaceborne or airborne optical instruments. All
548 simulations were carried out in the spectral range from 300 to 750 nm using 5-nm spectral bands and the results
549 were produced for the nominal wavelengths of each of the 81 bands, that is at 350, 355, 360, etc..., 745, 750 nm.
550 The results in the 300–350 nm range were not retained in the output files (that include seawater IOPs, radiometric
551 quantities and AOPs) because this spectral region was included primarily to account for potential effects of
552 inelastic processes at wavelengths longer than 350 nm and, additionally, it is known the uncertainties in the
553 characterization of seawater IOPs can increase significantly at wavelengths shorter than 350 nm.

554 For 3320 scenarios of input IOPs we performed several separate sets of RT simulations ~~which~~that differed in
555 terms of assumed sea-surface boundary conditions and the inclusion or exclusion of inelastic radiative processes
556 within the water column. The assumptions regarding the sea-surface boundary conditions were the same as in the
557 previous RT simulations described in Loisel et al. (2018). Specifically, all simulations were made under the same
558 assumption of wind speed of 5 m s^{-1} , which determines the sea-surface roughness involved in the calculations of
559 transmission and reflection of light at the air-water interface. In all simulations the sky conditions were also
560 assumed to be the same, i.e., clear skies and standard atmosphere. However, three distinct sets of simulations were
561 made for the three values of sun zenith angle, 0° , 30° , and 60° . With regard to consideration of inelastic processes,
562 we also performed three distinct sets of simulations. The first of these sets assumed the absence of inelastic
563 processes in water, that is no Raman scattering by water molecules, no fluorescence by chlorophyll-a, and no
564 fluorescence by CDOM. The second set of these simulations included Raman scattering by water molecules.
565 Finally, the third set included both Raman scattering and chlorophyll-a fluorescence, and this scenario of inelastic
566 processes is expected to generally provide the most realistic simulations of radiative transfer in the ocean surface
567 layer. We note, however, that fluorescence by CDOM was not included in any simulations. The Raman scattering
568 coefficient, phase function, and wavelength distribution function were set to their default values described in
569 HydroLight technical documentation (Mobley, 2012). The quantum efficiency of chlorophyll-a fluorescence,
570 [which may exhibit significant -variability \(nearly 5-fold between about 0.01 and 0.05\) in ocean waters \(Maritorena](#)
571 [et al., 2000; Morrison et al., 2003\)](#), was also set to its default value of 0.02 in the HydroLight code. For each
572 scenario of sun zenith angle and inelastic processes, we performed 3320 RT simulations, each for a different
573 combination of seawater IOPs. Thus, given the three sun zenith angles, the three scenarios of inelastic processes,
574 and 3320 combinations of IOPs, overall we performed 29880 simulations. The combination of the synthetic IOP
575 dataset used as input to RT simulations (section 2) and the results for the radiance, other radiometric quantities,
576 and AOPs obtained from these 29880 simulations (described in this section) constitute the synthetic ocean optical
577 database developed in this study.

578 **4 Comparisons of the synthetic database with in situ data**

579 In this section we compare the selected spectral IOP coefficients from the synthetic IOP dataset with in situ
580 data of IOPs and the selected spectral AOPs from the synthetic database generated with the RT simulations with

581 in situ data of AOPs. In these comparisons, we also include some empirical relationships between the IOPs or
 582 AOPs ~~which that~~ were established in previous studies based on the analysis of in situ data.
 583



584
 585 Figure 7. (a) $a_g(443)$, (b) $a_{dg}(443)$, (c) $a_d(443)$, and (d) $b_{bp}(550)$ as a function of $a_{ph}(443)$ for the in situ dataset (black data
 586 points) and the synthetic dataset (colored data points). The black polygon lines in each panel delimit approximately the scatter
 587 of the in situ data points (black dots). Each color refers to the optical water group as indicated (139, 262, and 2919 data points
 588 for Group 1, 2, and 3, respectively). Empirical relationships previously developed for (a) $a_g(443)$ vs. $a_{ph}(443)$ and (d) $b_{bp}(550)$
 589 vs. $a_{ph}(443)$ are also displayed for comparison. The original relationships were formulated as a function of Chla and the
 590 presented relationships were obtained by converting Chla to $a_{ph}(443)$ using the chlorophyll-specific phytoplankton absorption
 591 at 443 nm from Bricaud et al. (1998).
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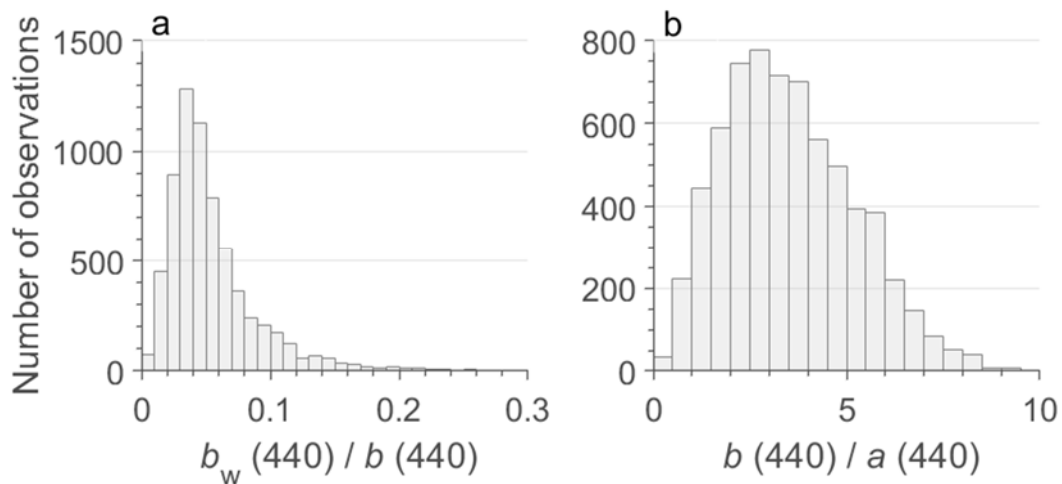
593 Figure 7 depicts the scatter plots of IOP coefficients, specifically $a_g(440)$ vs. $a_{ph}(440)$ (Fig. 7a), $a_{dg}(440)$ vs.
 594 $a_{ph}(440)$ (Fig. 7b), $a_d(440)$ vs. $a_{ph}(440)$ (Fig. 7c), and $b_{bp}(550)$ vs. $a_{ph}(440)$ (Fig. 7d). The scatter plots include two
 595 datasets, the in situ dataset and the synthetic dataset as described in section 2. We recall that in both types of
 596 datasets, $a_{ph}(440)$ plotted on the x -axis is the same because the phytoplankton absorption data used in this study
 597 were obtained from field measurements with no modeling involved. The scatter plots show a significant degree of
 598 overlap which indicates general consistency between the synthetic and in situ datasets. Similar patterns are
 599 observed when the $a_g(440)/a_{ph}(440)$, $a_{dg}(440)/a_{ph}(440)$, $a_d(440)/a_{ph}(440)$, and $b_{bp}(550)/a_{ph}(440)$ ratios are plotted
 600 as a function of $a_{ph}(440)$ (not shown). For illustrative purposes, the data from the synthetic IOP dataset are color
 601 coded to indicate the partitioning of data into the three OWGs, i.e., Groups 1, 2, and 3 that were defined using the
 602 synthetic spectra of $R_{rs}(\lambda)$ generated through RT simulations with input of the synthetic IOP data. As expected,
 603 the data with generally lowest values of IOPs belong to Group 3, the data with intermediate values of IOPs to
 604 Group 2, and the data with the highest IOPs (most turbid waters) to Group 1. We also note the in situ dataset
 605 exhibits somewhat wider dynamic range of variability than the synthetic dataset, especially when the IOP ratios,
 606 $a_g(440)/a_{ph}(440)$, $a_{dg}(440)/a_{ph}(440)$, $a_d(440)/a_{ph}(440)$, and $b_{bp}(550)/a_{ph}(440)$, are relatively high. While this result
 607 can reflect some degree of intrinsic difference in the dynamic range covered by the two datasets, it must also be
 608 recognized that some variability in the in situ dataset may be associated with the fact that these data were collected

609 on numerous cruises by different groups of investigators using the methodology (instrumentation, data processing,
610 data quality control, etc.) that unavoidably was not the same across all different data sources.

611 For additional comparative purposes, Fig. 7a,de includes a few empirical relationships between the IOPs in
612 question, which were established in previous studies based on considerable amount of field measurements
613 collected mostly in open ocean environments. As seen, the relationships between $a_g(440)$ and $a_{ph}(440)$ based on
614 the studies of Morel (2009) and Bricaud et al. (2010) agree quite well with the central tendency of variation within
615 our synthetic dataset. We note that Morel (2009) and Bricaud et al. (2010) reported on the relationships between
616 $a_g(440)$ and Chla ~~which that~~ were very similar in these studies. For the purpose of illustration in our Fig. 7a, we
617 replaced Chla with $a_{ph}(440)$ using the formula $Chla = a_{ph}(440)/0.05582$. Similarly, the studies of Huot et al. (2008)
618 and Antoine et al. (2011) reported on empirical relationships between $b_{bp}(\lambda)$ and Chla. After converting Chla to
619 $a_{ph}(440)$ as mentioned above, these two relationship are plotted in Fig. 7de. Although these two relationships have
620 different slopes, they are both generally consistent with the average trend of variation in the synthetic dataset.

621 The radiative transfer is driven mainly by two main ratios of IOPs which are the scattering to absorption ratio,
622 $b(\lambda)/a(\lambda)$, which controls the number of scattering events (Morel and Gentili, 1991), and the molecular to total
623 scattering ratio, $b_w(\lambda)/b(\lambda)$, which is the parameter controlling the weighted sum of the particle- scattering and
624 molecular scattering phase functions (Morel and Loisel, 1998; Loisel and Stramski, 2000). Figure 8 shows the
625 distribution of these two ratios at 440 nm for the synthetic data set dataset. The $b_w(440)/b(440)$ and $b(440)/a(440)$
626 ratios range between about 0 and 0.2 and 0.5 and 10, respectively, which is consistent with previous models
627 developed for Case-1 waters (Figs. 2 and 3 in Morel and Gentili, 1991; Fig. 2 in Morel and Loisel, 1998;).

628



629

630 Figure 8. (a) Histograms of (a) $b_w(440)/b(440)$ and (b) $b(440)/a(440)$ for the synthetic dataset.

631

632

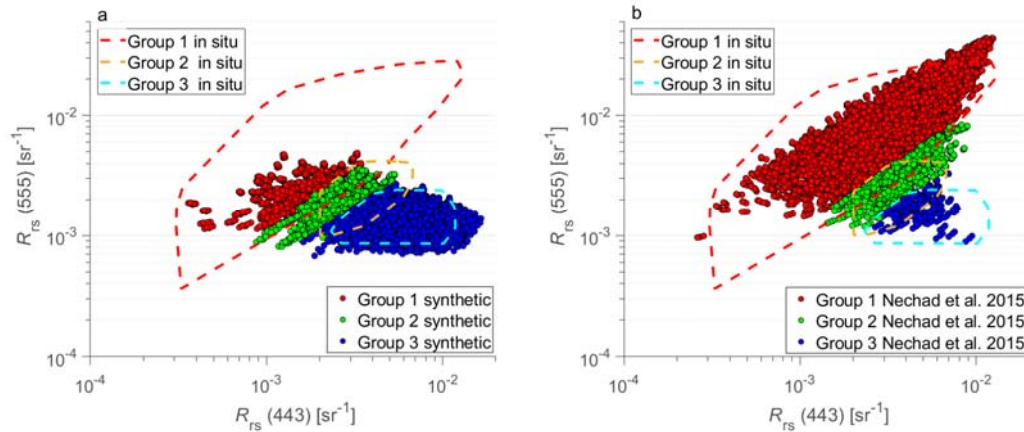
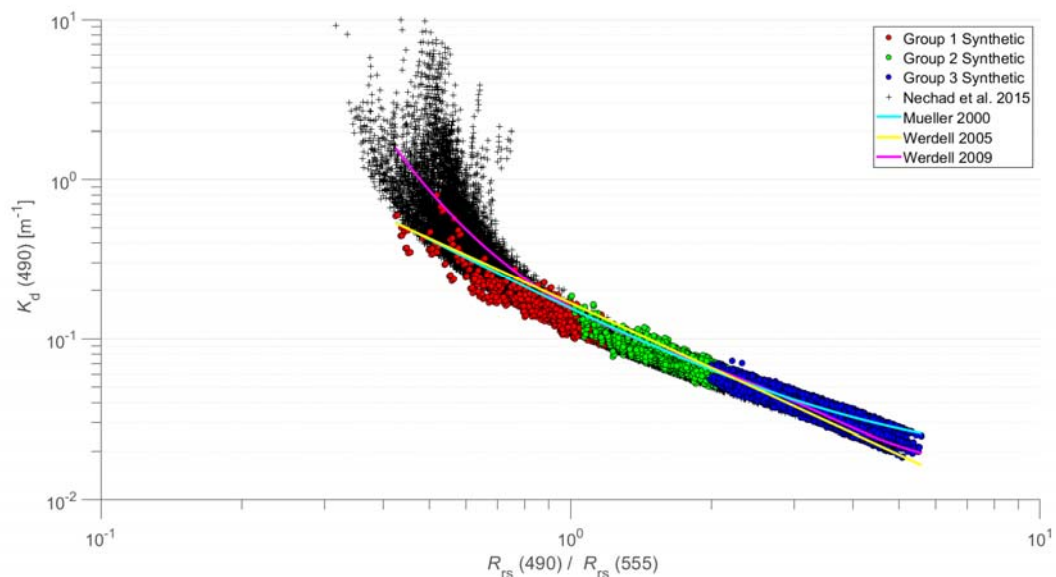


Figure 98. (a) $R_{rs}(555)$ as a function of $R_{rs}(443)$ for the synthetic dataset (colored data points) and in situ dataset (colored contours). (b) Same as panel (a) but for the synthetic dataset of Nechad et al. (2015) which that was developed for coastal waters. The color coding refers to the optical water groups as indicated.

For comparing the AOPs from the synthetic database with in situ data, we have chosen two AOPs, the spectral remote-sensing reflectance, $R_{rs}(\lambda)$, and the spectral diffuse attenuation coefficient of downwelling plane irradiance averaged within the water column from the sea surface to the first attenuation depth, $\langle K_d(\lambda) \rangle_1$, and the maximum band ratio of reflectance, MBR. The scatter plot of our synthetic data of $R_{rs}(555)$ vs. $R_{rs}(443)$ are depicted in Fig. 98a. For comparison, the range of in situ data is illustrated by the dashed contour lines. The maximum values of $R_{rs}(443)$ reached 0.0165 sr^{-1} , which is in good agreement with in situ measurements performed in ultra-oligotrophic waters in the South Pacific gyre during the BIOSOPE cruise (see Fig. 3 in Stramski et al., 2008). These results are once again illustrated using color coding to represent different optical water types, specifically Groups 1, 2, and 3. As seen, there is a relatively good agreement between the synthetic data and the range of variability of the in situ data for Groups 2 and 3 (Fig. 98a). For Group 1 (very turbid waters), however, the synthetic data exhibit a smaller range of variability compared with in situ data. This result is not unexpected because our primary goal was to generate the synthetic database that is most representative of open ocean pelagic environments as well as coastal areas where water turbidity is low to moderate rather than very high. As described in section 2, turbid waters of Group 1 correspond to Optical Water Classes 1 and 2 as defined in Mélin and Vantrepotte (2015). It is interesting to note that the synthetic optical database that was developed by Nechad et al. (2015) for coastal waters shows a relatively good consistency between the synthetic and in situ data for Group 1 (Fig. 8b). However, in contrast to our synthetic database, the synthetic data of Nechad et al. (2015) exhibit a limited range of variability compared with in situ data for Groups 2 and 3. Thus, the synthetic data of Nechad et al. (2015) for turbid waters in Group 1 can provide useful complementarity to our synthetic database whose main focus is on water types from Groups 2 and 3.

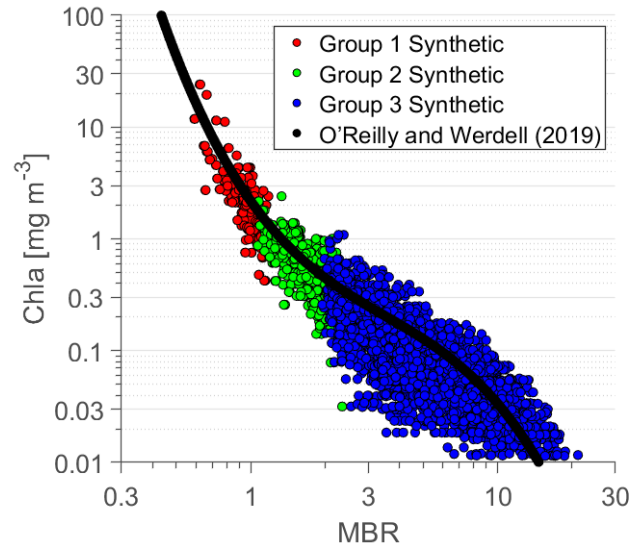
The scatter plot of the synthetic data of $\langle K_d(490) \rangle_1$ as a function of blue-to-green band ratio of reflectance, $R_{rs}(490)/R_{rs}(555)$, is shown in Fig. 109. These synthetic data are again color coded according to optical water classes defined in terms of Groups, 1, 2, and 3. For comparison, a few empirical relationships between these AOP variables established in previous analyses of field measurements are also displayed in Fig. 109 (Mueller, 2000; Werdell, 2005; Werdell, 2009). The relationship of Mueller (2000) was formulated during the early phase of SeaWiFS satellite mission to serve as an operational global algorithm for estimating $K_d(490)$ from ocean color observations. Werdell (2005) provided an updated relationship with a primary goal to improve the estimation of

666 $K_d(490)$ at low values of $K_d(490)$ that correspond to high values of reflectance band ratio. Figure 109 shows that
 667 these two relationships are generally consistent with our synthetic data across the entire range of variability
 668 encompassing data from Groups, 1, 2, and 3. This is reassuring given that the main purpose of our synthetic
 669 database and these two empirical relationships is similar in a sense of targeting the optical variability within the
 670 global ocean dominated by open ocean environments. Figure 109 also includes the relationship of Werdell (2009)
 671 ~~which that~~ represents the most recent update of global empirical algorithms for estimating $K_d(490)$ from different
 672 ocean color satellite sensors. Specifically, the relationship of Werdell (2009) presented in Fig. 9 is referred to as
 673 KD2S and is based on SeaWiFS spectral bands. In contrast to relationships of Mueller (2020) and Werdell (2005),
 674 the relationship of Werdell (2009) deviates significantly from our synthetic data within the range of relatively high
 675 values of $\langle K_d(490) \rangle_1$ which correspond to relatively low values of $R_{rs}(490)/R_{rs}(555)$. It is remarkable that this
 676 deviation occurs within the range where our synthetic data are classified as Group 1, so these are the optical water
 677 types associated with high water turbidity. Another remarkable result illustrated in Fig. 109 is that the relationship
 678 of Werdell (2009) in this range is quite consistent with the main trend observed within the synthetic database of
 679 Nechad et al. (2015) ~~which that~~ was developed for coastal environments. This result further supports the potential
 680 complementarity between our synthetic database and database of Nechad et al. (2015).
 681



682
 683 Figure 10. Scatter plot of $K_d(490)$ vs. the blue-to-green reflectance ratio, $R_{rs}(490)/R_{rs}(555)$, for the synthetic database. The red,
 684 green, and blue data points represent the three optical water groups 1, 2, and 3, respectively. The black cross-mark data points
 685 are from the Nechad et al. (2015) synthetic dataset. The curves representing the relationships developed by Mueller (2020),
 686 Werdell (2005), and Werdell (2009) are also displayed. The $K_d(490)$ data points represent $\langle K_d(490) \rangle_1$ for the present synthetic
 687 database (colored data points), and the near-surface $K_d(490)$ calculated within the top 1 cm layer for the Nechad dataset (black
 688 data points).
 689

690 The scatter plot of Chla vs. the maximum band ratio of reflectance, MBR, for the synthetic database is shown
 691 in Fig. 11. The monotonically decreasing trend of Chla with increasing MBR is consistent with the SeaWiFS-
 692 specific OC5 algorithm for estimating Chla from MBR (O'Reilly and Werdell, 2019). For this illustration, we
 693 estimated Chla using the relationship between $a_{ph}(660)$ and Chla from Bricaud et al. (1998), which is unavoidably
 694 affected to some extent by natural variability in this relationship.



695

696 [Figure 11. Scatter plot of Chla vs. the blue-to-green maximum band ratio \(MBR\) of remote-sensing reflectance \(i.e. \$R_{rs}\(412 > 443 > 490 > 510\)/R_{rs}\(555\)\$ \) for the synthetic database. The red, green, and blue data points represent the three optical](#)
 697 [water groups 1, 2, and 3, respectively. The solid black line represents the OC5 algorithm developed by O'Reilly and Werdell](#)
 698 [\(2019\) for SeaWiFS spectral bands. For this illustration, Chla was calculated from \$a_{ph}\(660\)\$ using the chlorophyll-specific](#)
 699 [phytoplankton absorption at 660 nm from Bricaud et al. \(1998\).](#)
 700

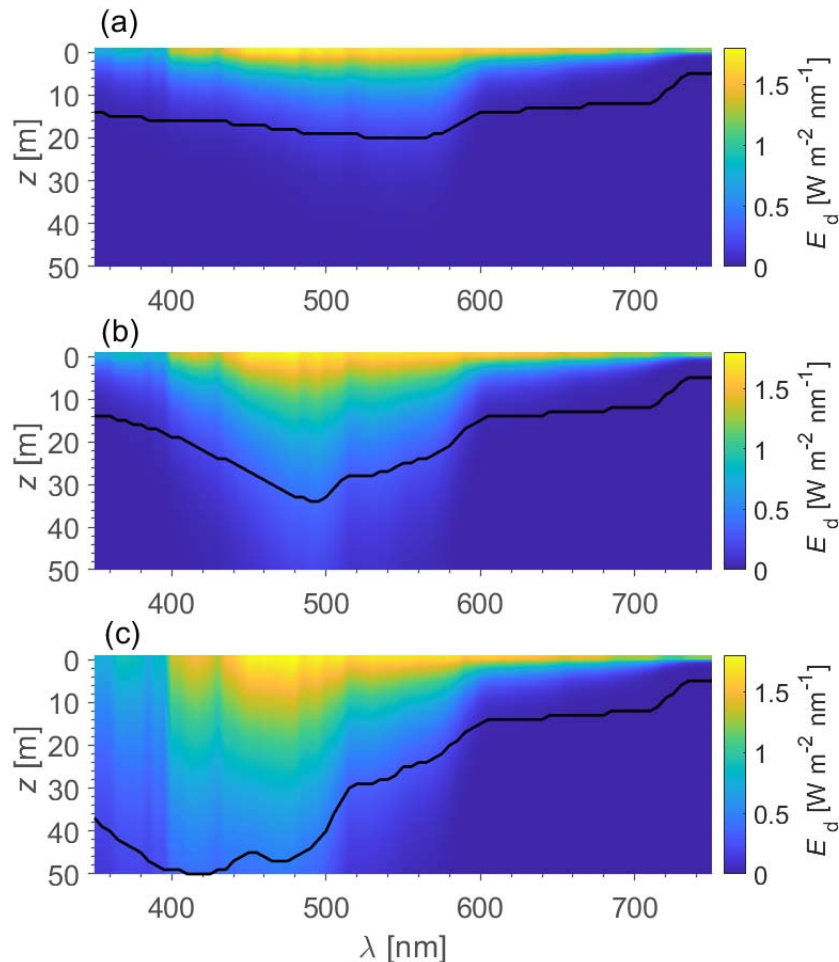
701

702 5 Summary

703 We have generated a new synthetic database that consists of seawater IOPs as well as corresponding
 704 radiometric quantities and AOPs within the ocean surface layer down to a depth of 50 m and at the sea surface.
 705 The radiometric quantities and AOPs were obtained from radiative transfer (RT) simulations performed with
 706 Hydrolight code using the IOPs as input to the calculations. The list of variables included in the database is
 707 provided in Table 2. [Because of the use of absorption and scattering properties of pure seawater \(assuming the](#)
 708 [salinity of 35‰\) in the simulations, the present database cannot be used for applications to freshwater](#)
 709 [environments and also special caution should be exercised for applications when water salinity is significantly less](#)
 710 [than 35‰ because of decrease in pure seawater scattering.](#) This database is organized following an easy to read
 711 netcdf structure and divided into two subsets of data for which the file name identifies the sun zenith angle and the
 712 RT simulation scenario related to the presence or absence of inelastic radiative processes within the water column.
 713 The first subset of data includes the seawater spectral absorption and backscattering coefficients as well as sea-
 714 surface radiometric quantities relevant to ocean color radiometry, $R_{rs}(\lambda)$, $L_w(\lambda)$, $E_d(z=0^+, \lambda)$, and $L_u(z=0^+, \lambda)$ where
 715 $z=0^+$ is just above the surface. The surface and depth-profile values of several spectral radiometric quantities and
 716 AOPs, as well as PAR are included in the second subset of data. The spectra of z_{eu} and z_1 are also provided in the
 717 second file. More details on the organization and content of the database are included in readme file that is also
 718 provided in the database.

719 In closing, we present an example illustration of one of the radiometric variables included in the output data
 720 files generated by RT simulations. We recall that the primary result of HydroLight simulations is the spectral
 721 radiance that provides a comprehensive information about the angular distribution of light field, from which
 722 different irradiances and AOPs are calculated. However, it is the spectral downwelling plane irradiance, $E_d(z, \lambda)$,
 723 [which that](#) has been the most commonly measured radiometric quantity in ocean optics, so in Fig. 12~~0~~ we have

724 chosen to illustrate the HydroLight-simulated $E_d(z, \lambda)$ within the ocean surface layer down to a depth of 50 m.
 725 These results are presented for three different scenarios of IOPs which are representative of three different optical
 726 water types defined in terms of Group 1, Group 2, and Group 3 (see section 2). These RT simulations were
 727 performed for the sun zenith angle of 30° in the presence of Raman scattering by water molecules and chlorophyll-
 728 a fluorescence in the water column. In addition to significant differences in the variation of the spectral $E_d(\lambda)$ as a
 729 function of depth z between the Groups 1, 2, and 3, Fig. 120 also illustrates distinct differences in the magnitude
 730 and spectral behavior of the first optical attenuation depth, z_1 . This quantity is equivalent to the inverse of the
 731 diffuse attenuation coefficient, $\langle K_d(\lambda) \rangle_1$. As expected, the first attenuation depth z_1 is located much closer to the
 732 ocean surface for data from Group 1 (Fig. 120a) compared with Group 2 (Fig. 120b) and Group 3 (Fig. 120c),
 733 especially across the blue-green region of the spectrum. In the red part of the spectrum where pure water absorption
 734 dominates the attenuation of $E_d(\lambda)$, the differences between the three groups are small. It is also notable that the
 735 spectral behavior of z_1 for Group 3 (Fig. 120c) that represents relatively clear ocean waters is remarkably similar
 736 to the spectral shape of pure water absorption coefficient.
 737



738
 739 Figure 120. Examples of depth profiles of $E_d(z, \lambda)$ for a given IOP scenario from (a) the optical water group (OWG) 1, (b)
 740 OWG 2, and (c) OWG 3. Radiative transfer simulations were performed for a sun zenith angle of 30° and included Raman
 741 scattering by water molecules and chlorophyll-a fluorescence.

742

744 **Table 2:** Symbols, variables, and units for the various quantities included in the final synthetic optical database.

Symbol	Variable*	Units
z	Depth in water	m
λ	Light wavelength in vacuum	nm
a, b, b_b	Total absorption, scattering, and backscattering coefficients of seawater	m^{-1}
a_{nw}	Absorption coefficient of all non-water constituents	m^{-1}
a_{ph}, a_d, a_g	Absorption coefficients of phytoplankton, non-algal particles, and CDOM	m^{-1}
b_{nw}	Backscattering coefficient of all non-water constituents	m^{-1}
b_{b-ph}, b_{b-d}	Backscattering coefficients of phytoplankton and non-algal particles	m^{-1}
b_{nw}	Scattering coefficient of all non-water constituents	m^{-1}
b_{ph}, b_d	Scattering coefficients of phytoplankton and non-algal particles	m^{-1}
E_o, E_{od}, E_{ou}	Total, downwelling, and upwelling scalar irradiances	$W m^{-2} nm^{-1}$
E_d, E_u	Downwelling and upwelling plane irradiances	$W m^{-2} nm^{-1}$
L_w, L_u	Water-leaving and upwelling radiances	$W m^{-2} sr^{-1} nm^{-1}$
PAR	Photosynthetically Available Radiation defined as the total quantum scalar irradiance within the spectral range 400-700 nm	$\mu mol photons s^{-1} m^{-2}$
R_{rs}	Remote-sensing reflectance	sr^{-1}
K_x	Diffuse attenuation coefficients for upwelling and downwelling plane irradiances or upwelling radiance (the radiometric quantity is indicated by subscript x)	m^{-1}
μ_d, μ_u	Average cosines of downwelling and upwelling light fields	dimensionless
z_{eu}	Euphotic depth at which PAR is reduced to 1% of its surface value	m
z_1	First optical attenuation depth at which spectral E_d or PAR is reduced to 36.8% of its surface value	m

745 *All optical variables in the database are spectral and provided at different light wavelengths between 350 and
 746 750 nm at 5 nm intervals and different depths within the water column between the sea surface and the 50 m
 747 depth, except for R_{rs} , and L_w which are defined at the sea surface.

748
 749 **Author contributions.** The concept of this study originated from the authors' discussions about the need for a
 750 new synthetic optical database in support of ocean color science and applications, especially the global ocean
 751 applications, including support of upcoming NASA's PACE hyperspectral ocean color satellite mission. All co-
 752 authors contributed to curation of in situ data. HL and DSFJ led the generation of the synthetic IOP dataset and
 753 created the satellite IOP dataset. DSFJ ran the RT simulations. HL and DS wrote the manuscript. All co-authors
 754 contributed to discussion, review, and editing of the manuscript.

755 **Competing interests.** The authors declare that they have no conflict of interest.

756 **Disclaimer.** Mention of trade names or commercial products does not constitute endorsement or recommendation
 757 for use. The views expressed in this article are those of the authors.

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762 **Data Availability:** The DOI (doi:10.6076/D1630T) is not yet active but reserved at the Dryad data repository. The
 763 database is available at Dryad open-access repository of research data (Loisel et al., 2023). Following completion
 764 of the review process,
 765 the synthetic optical database described in this study will be
 766 publicly available at the Dryad open-access repository of
 767 research data (Loisel et al., 2023; <https://doi.org/10.6076/D1630T>). ~~For the review purpose these data can be~~
 768 ~~downloaded using the following link:~~

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772 **Review statement.**

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